

# The ZIP Code of MSSM Particles

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# Where do MSSM fields live?

## Localization properties of quarks, leptons and Higgses

- Higgs bosons and top-quark in the “bulk” lead to large top-quark Yukawa coupling
- first 2 families localized (exhibiting family symmetries)

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## The legacy of higher dimensions

- Mirage Mediation (compressed spectrum for gauginos)
- Natural Susy
- discrete (nonabelian) family symmetries

Remnants of  $N=4$  SUSY from higher dimensions that might hide Susy at the LHC!

# Guidelines

- **Spinors if  $SO(10)$**  might be important even in absence of GUT gauge group
- gauge-top Yukawa unification in the MSSM
- presence of **discrete symmetries** with many applications

(Kobayashi, HPN, Ploeger, Raby, Ratz, 2006)

# Guidelines

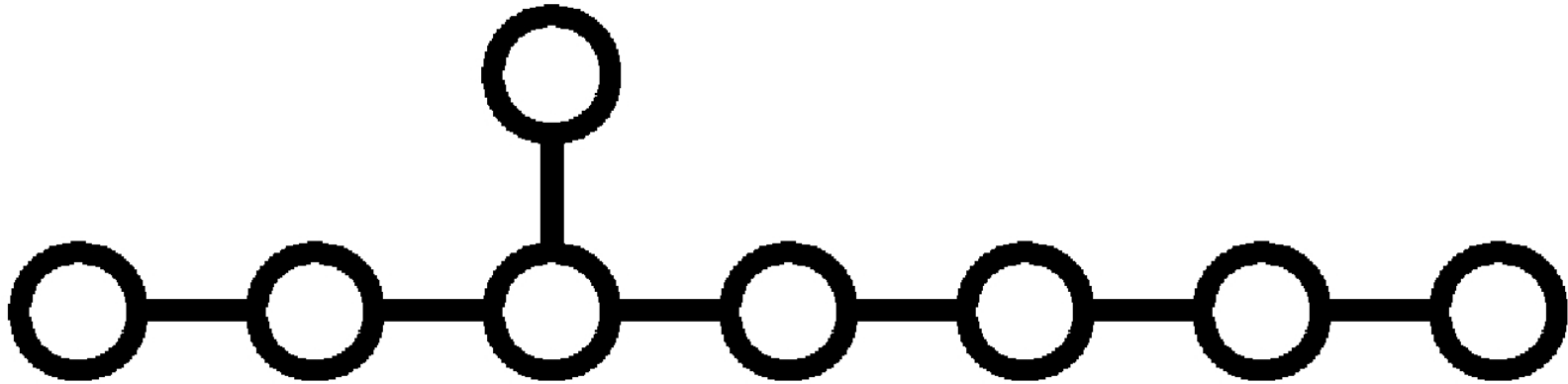
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(Kobayashi, HPN, Ploeger, Raby, Ratz, 2006)

From the mathematical structure we might prefer exceptional groups

- There is a maximal group:  $E_8$ ,
- but  $E_8$  and  $E_7$  do not allow chiral fermions in  $d = 4$ .
- How does this fit with our usual picture of unification based on  $SU(5)$  or  $SO(10)$ ?

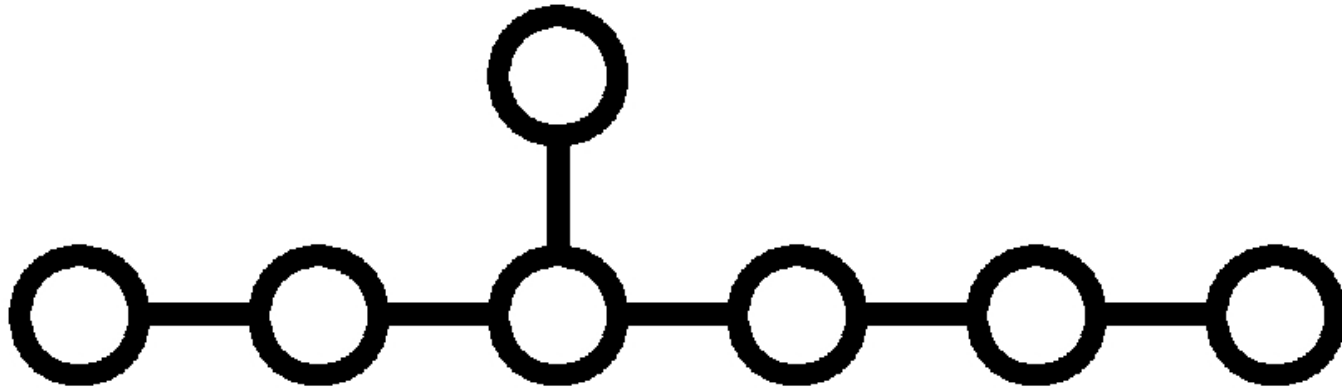
# Maximal Group



$E_8$  is the maximal group.

There are, however, no chiral representations in  $d = 4$ .

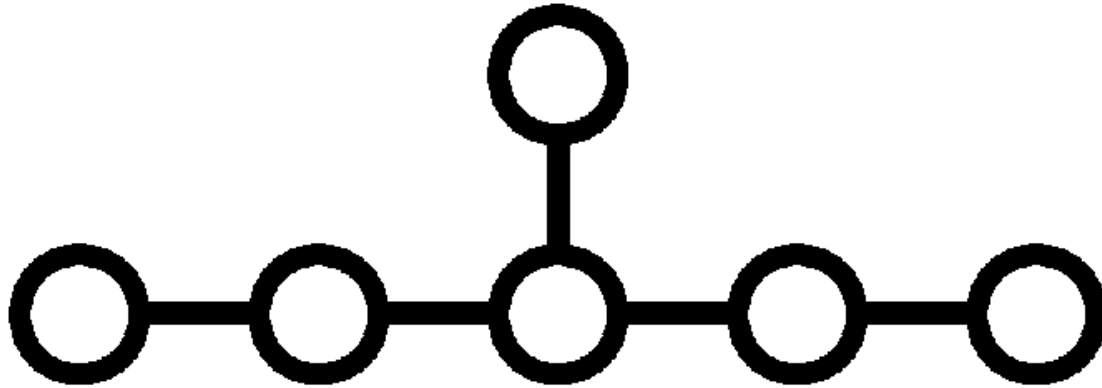
$E_7$



Next smaller is  $E_7$ .

No chiral representations in  $d = 4$  either

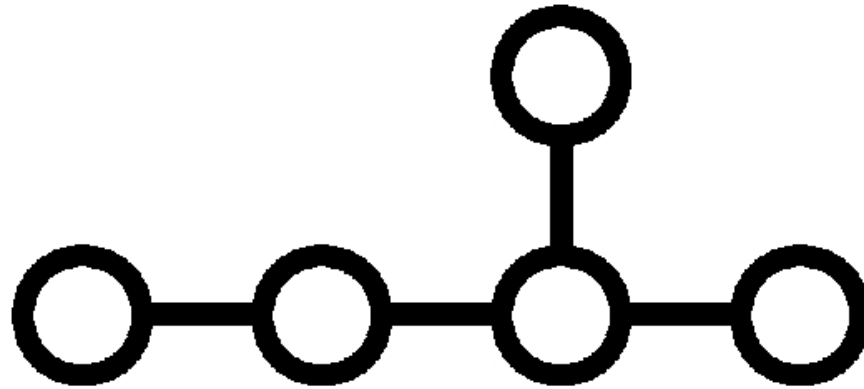
$E_6$



$E_6$  allows for chiral representations even in  $d = 4$ .



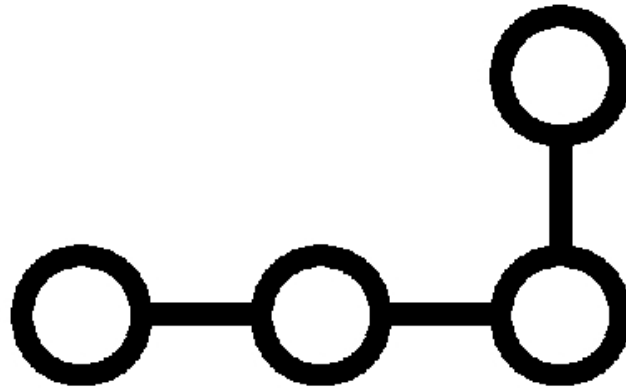
$$E_5 = D_5$$



$E_5$  is usually not called exceptional.

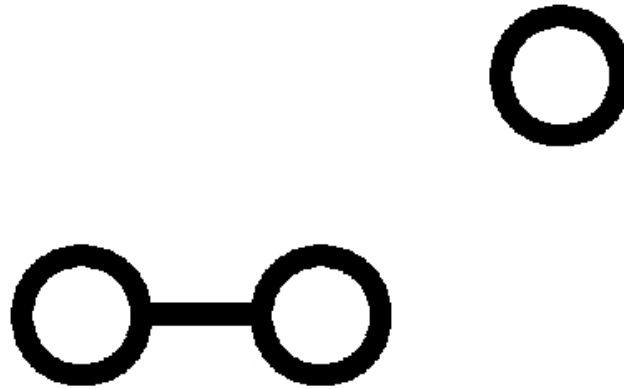
It coincides with  $D_5 = SO(10)$ .

$$E_4 = A_4$$



$E_4$  coincides with  $A_4 = SU(5)$

$E_3$



$E_3$  coincides with  $A_2 \times A_1$  which is  $SU(3) \times SU(2)$ .

# Exceptional groups in string theory

String theory “favours”  $E_8$

- $E_8 \times E_8$  heterotic string
- $E_8$  enhancement as a nonperturbative effect (M- or F-theory)

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Strings live in higher dimensions:

- chiral spectrum possible even with  $E_8$
- $E_8$  broken in process of compactification
- provides source for (nonabelian) discrete symmetries
- from  $E_8/SO(10)$  and  $SO(6)$  of the higher dimensional Lorentz group

# Geography

Many properties of the models depend on the geography of extra dimensions, such as

- the **location** of quarks and leptons,
- the **relative location** of Higgs bosons,

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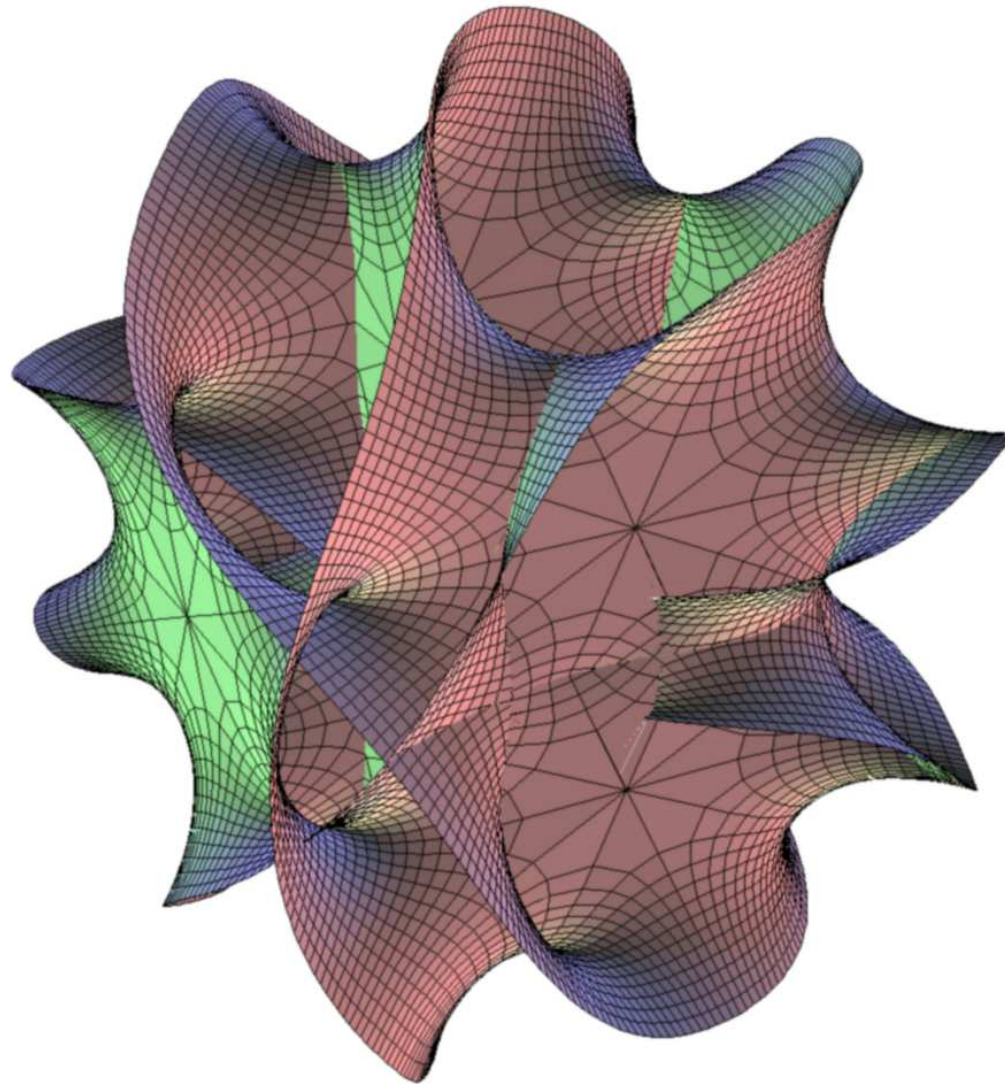
- the **location** of quarks and leptons,
- the **relative location** of Higgs bosons,

but there is also a “localization” of gauge fields

- $E_8 \times E_8$  in the bulk
- smaller gauge groups on various branes

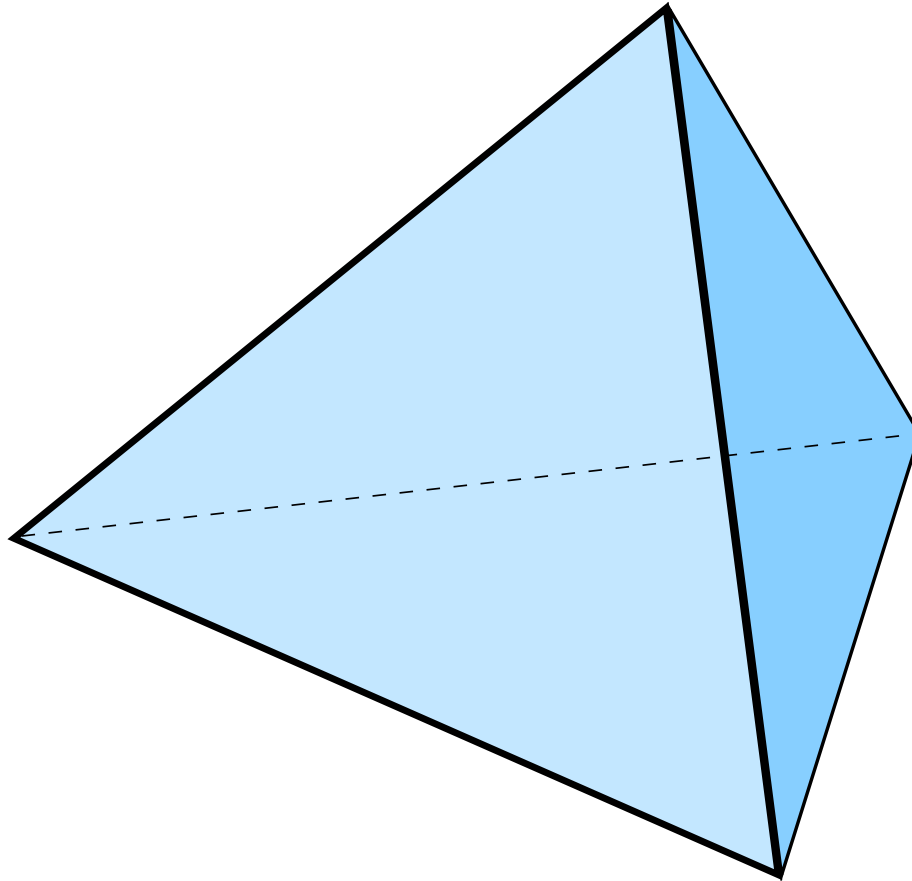
Observed 4-dimensional gauge group is common subgroup of the various localized gauge groups!

# Calabi Yau Manifold





# Orbifold



(Dixon, Harvey, Vafa, Witten, 1985)

# Localization

Quarks, Leptons and Higgs fields can be localized:

- in the Bulk ( $d = 10$  **untwisted** sector)
- on 3-Branes ( $d = 4$  twisted sector **fixed points**)
- on 5-Branes ( $d = 6$  twisted sector **fixed tori**)

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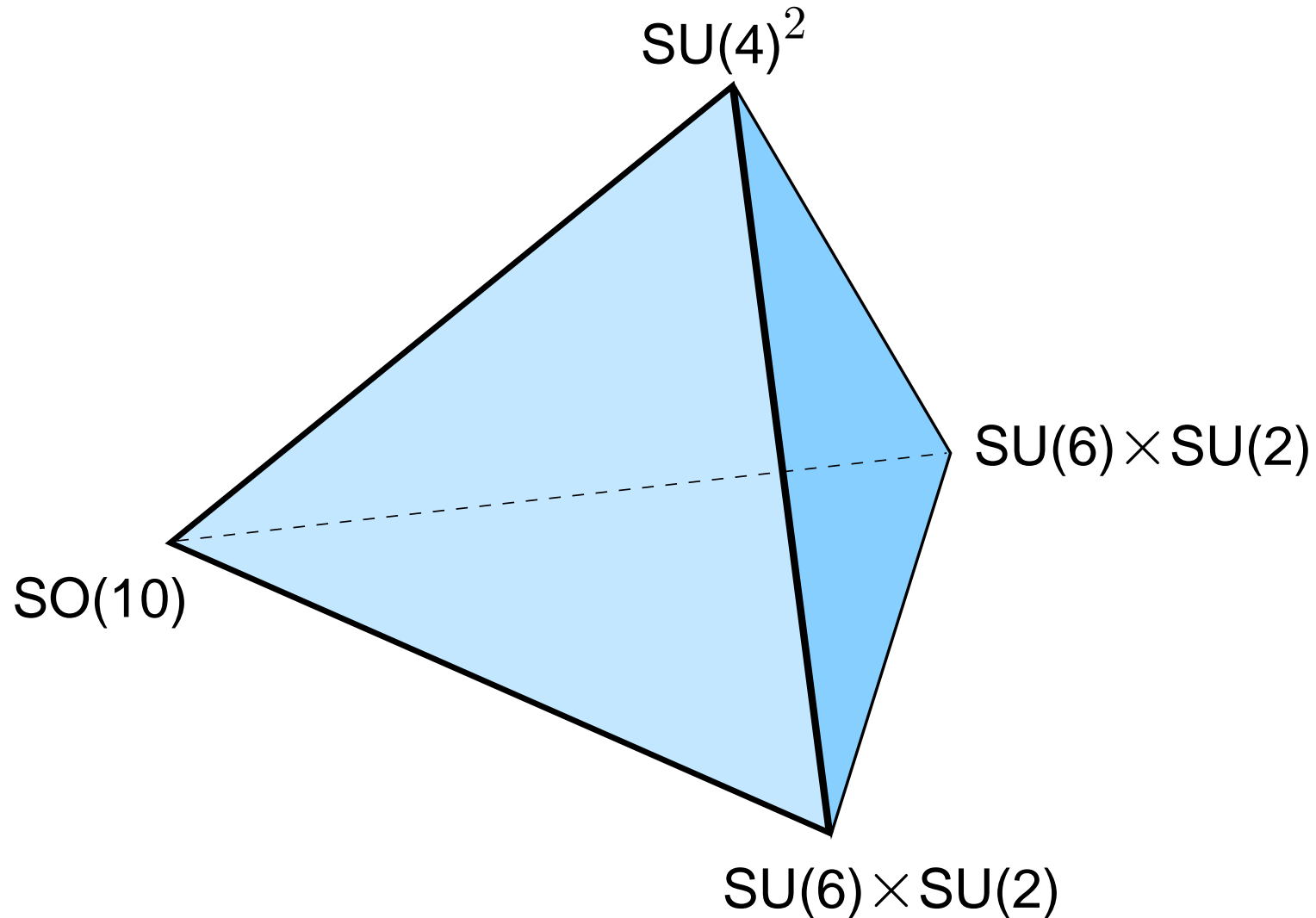
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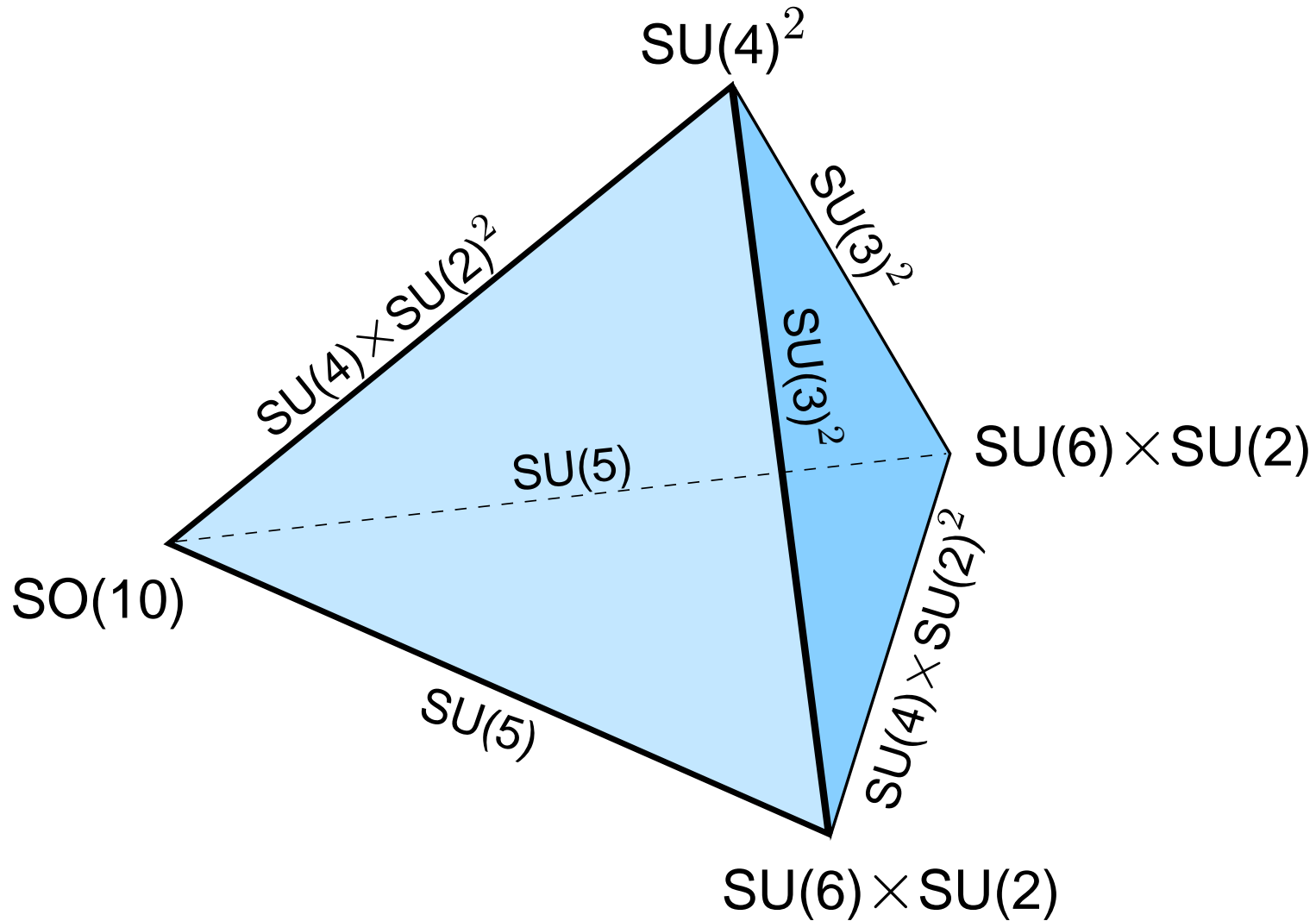
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# Localized gauge symmetries



(Förste, HPN, Vaudrevange, Wingerter, 2004)

# Standard Model Gauge Group



# The (extended) MiniLandscape

- in  $Z_6II$  orbifold many models with the exact spectrum of the MSSM (absence of chiral exotics)

(Lebedev, HPN, Raby, Ramos-Sanchez, Ratz, Vaudrevange, Wingerter, 2006-2009)

- recently extended to the  $Z_2 \times Z_4$  orbifold

(Mayorga Pena, Nilles, Oehlmann, 2012)

- large top quark Yukawa coupling

- family symmetries for the first two families

(Kobayashi, HPN, Ploeger, Raby, Ratz, 2006)

- models with R-parity + solution to the  $\mu$ -problem

(Lebedev, HPN, Raby, Ramos-Sanchez, Ratz, Vaudrevange, Wingerter, 2007)

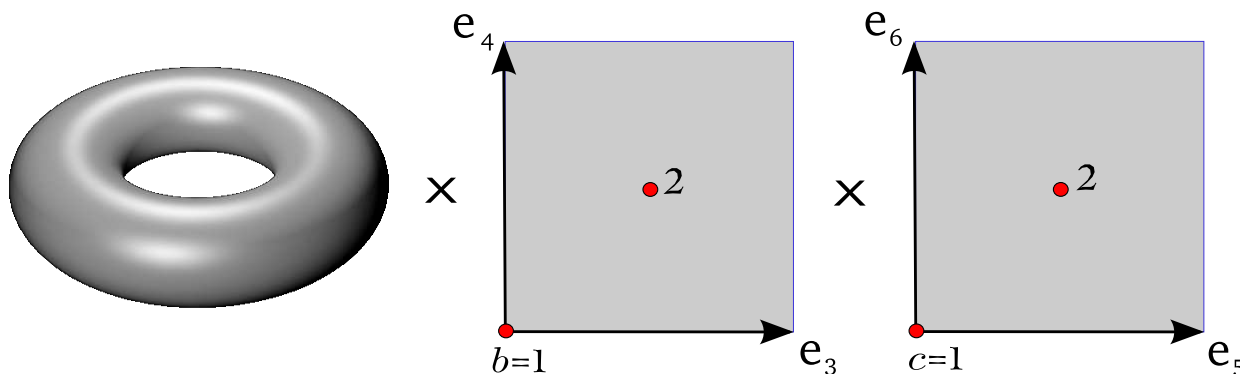
- gaugino condensation and mirage mediation

(Löwen, HPN, 2008)

# Sectors

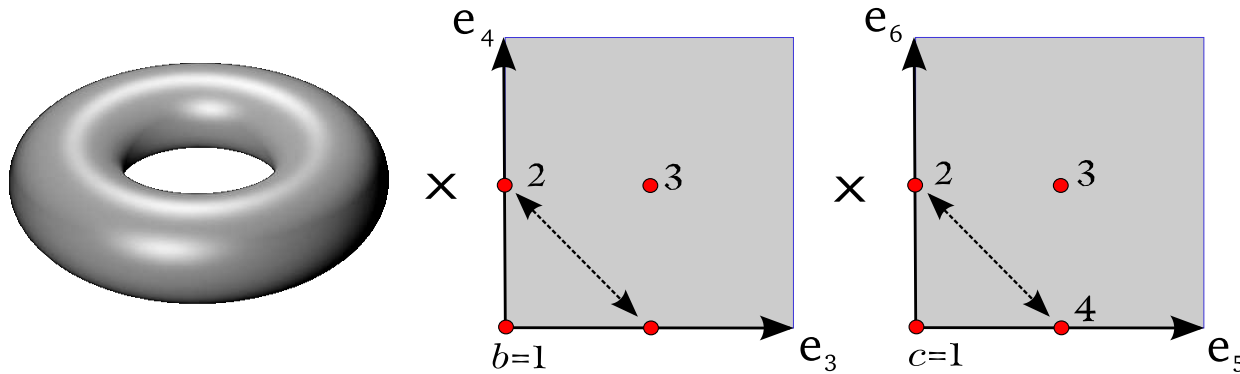
The underlying  $Z_2 \times Z_4$  orbifold has the following sectors:

- the **untwisted sector** ( bulk  $D = 10$ ,  $N = 4$  Susy)
- twisted sectors corresponding to  $Z_2$  ( $\theta$ ) and  $Z_4$  ( $\omega$ ) twists



The  $\omega$  sector has  $2 \times 2 = 4$  fixed tori, corresponding to “5-branes” confined to  $D=6$  space-time ( $N = 2$  Susy).

# $\omega^2$ twisted sector

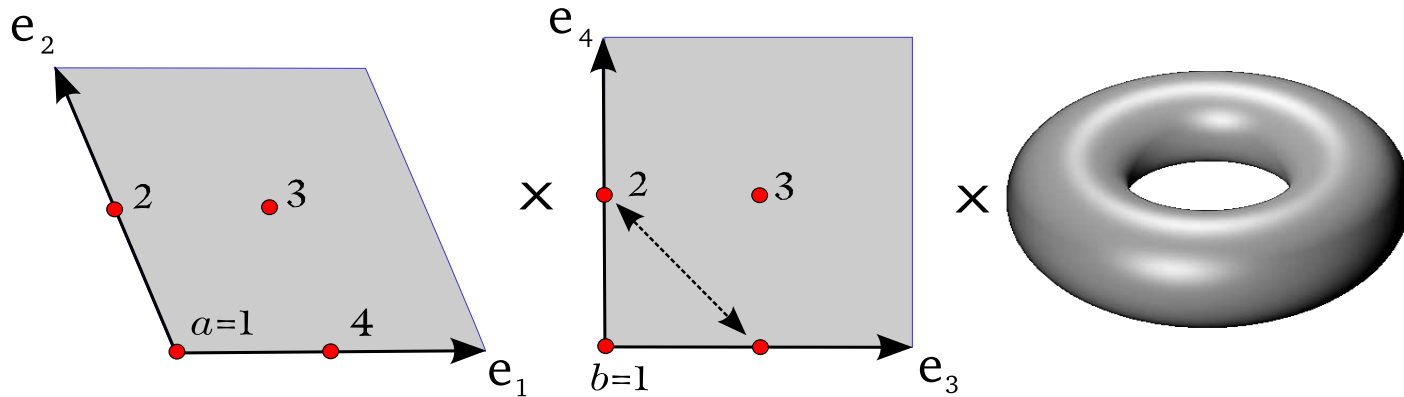


The  $\omega^2$  sector contains fixed tori corresponding to

- “5-branes” confined to 6 space-time dimensions (remnants of  $N = 2$  Susy)



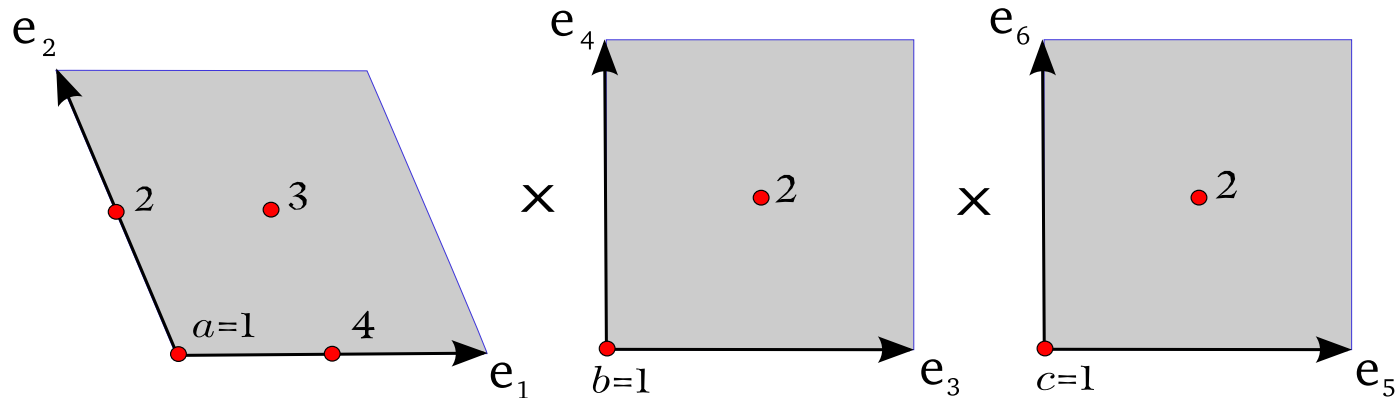
# $\theta$ twisted sector



The  $\theta$  sector contains  $4 \times 3$  fixed tori:

- “5-branes” confined to 6 space-time dimensions (sector with  $N = 2$  Susy)

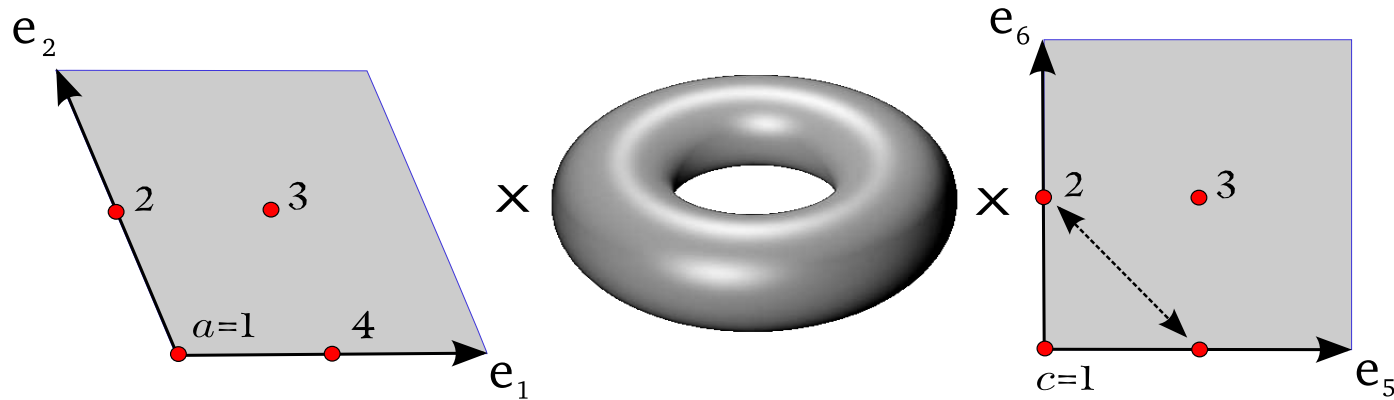
# $\theta\omega$ twisted sector



The  $\theta\omega$  sector contains  $4 \times 2 \times 2$  fixed points:

- “3-branes” confined to 4 space-time dimensions (sector with  $N = 1$  Susy)

# $\theta\omega^2$ twisted sector



The  $\theta\omega^2$  sector contains 4 x 3 fixed tori:

- “5-branes” confined to 6 space-time dimensions (sector with  $N = 2$  Susy)

Where do we find quarks, leptons and Higgs bosons in the models of the MiniLandscape?

# A Benchmark Model

At the orbifold point the gauge group is

$$SU(3) \times SU(2) \times U(1)^9 \times SU(4) \times SU(2)$$

- one  $U(1)$  is anomalous
- there are singlets and vectorlike exotics
- decoupling of exotics and breakdown of gauge group has been verified
- remaining gauge group

$$SU(3) \times SU(2) \times U(1)_Y \times SU(4)_{\text{hidden}}$$

- for discussion of neutrinos and R-parity we keep also the  $U(1)_{B-L}$  charges

# Spectrum

#	irrep	label	#	irrep	label
3	$(\mathbf{3}, \mathbf{2}; \mathbf{1}, \mathbf{1})_{(1/6, 1/3)}$	$q_i$	3	$(\bar{\mathbf{3}}, \mathbf{1}; \mathbf{1}, \mathbf{1})_{(-2/3, -1/3)}$	$\bar{u}_i$
3	$(\mathbf{1}, \mathbf{1}; \mathbf{1}, \mathbf{1})_{(1, 1)}$	$\bar{e}_i$	8	$(\mathbf{1}, \mathbf{2}; \mathbf{1}, \mathbf{1})_{(0, *)}$	$m_i$
3 + 1	$(\bar{\mathbf{3}}, \mathbf{1}; \mathbf{1}, \mathbf{1})_{(1/3, -1/3)}$	$\bar{d}_i$	1	$(\mathbf{3}, \mathbf{1}; \mathbf{1}, \mathbf{1})_{(-1/3, 1/3)}$	$d_i$
3 + 1	$(\mathbf{1}, \mathbf{2}; \mathbf{1}, \mathbf{1})_{(-1/2, -1)}$	$l_i$	1	$(\mathbf{1}, \mathbf{2}; \mathbf{1}, \mathbf{1})_{(1/2, 1)}$	$\bar{l}_i$
1	$(\mathbf{1}, \mathbf{2}; \mathbf{1}, \mathbf{1})_{(-1/2, 0)}$	$h_d$	1	$(\mathbf{1}, \mathbf{2}; \mathbf{1}, \mathbf{1})_{(1/2, 0)}$	$h_u$
6	$(\bar{\mathbf{3}}, \mathbf{1}; \mathbf{1}, \mathbf{1})_{(1/3, 2/3)}$	$\bar{\delta}_i$	6	$(\mathbf{3}, \mathbf{1}; \mathbf{1}, \mathbf{1})_{(-1/3, -2/3)}$	$\delta_i$
14	$(\mathbf{1}, \mathbf{1}; \mathbf{1}, \mathbf{1})_{(1/2, *)}$	$s_i^+$	14	$(\mathbf{1}, \mathbf{1}; \mathbf{1}, \mathbf{1})_{(-1/2, *)}$	$s_i^-$
16	$(\mathbf{1}, \mathbf{1}; \mathbf{1}, \mathbf{1})_{(0, 1)}$	$\bar{n}_i$	13	$(\mathbf{1}, \mathbf{1}; \mathbf{1}, \mathbf{1})_{(0, -1)}$	$n_i$
5	$(\mathbf{1}, \mathbf{1}; \mathbf{1}, \mathbf{2})_{(0, 1)}$	$\bar{\eta}_i$	5	$(\mathbf{1}, \mathbf{1}; \mathbf{1}, \mathbf{2})_{(0, -1)}$	$\eta_i$
10	$(\mathbf{1}, \mathbf{1}; \mathbf{1}, \mathbf{2})_{(0, 0)}$	$h_i$	2	$(\mathbf{1}, \mathbf{2}; \mathbf{1}, \mathbf{2})_{(0, 0)}$	$y_i$
6	$(\mathbf{1}, \mathbf{1}; \mathbf{4}, \mathbf{1})_{(0, *)}$	$f_i$	6	$(\mathbf{1}, \mathbf{1}; \bar{\mathbf{4}}, \mathbf{1})_{(0, *)}$	$\bar{f}_i$
2	$(\mathbf{1}, \mathbf{1}; \mathbf{4}, \mathbf{1})_{(-1/2, -1)}$	$f_i^-$	2	$(\mathbf{1}, \mathbf{1}; \bar{\mathbf{4}}, \mathbf{1})_{(1/2, 1)}$	$\bar{f}_i^+$
4	$(\mathbf{1}, \mathbf{1}; \mathbf{1}, \mathbf{1})_{(0, \pm 2)}$	$\chi_i$	32	$(\mathbf{1}, \mathbf{1}; \mathbf{1}, \mathbf{1})_{(0, 0)}$	$s_i^0$
2	$(\bar{\mathbf{3}}, \mathbf{1}; \mathbf{1}, \mathbf{1})_{(-1/6, 2/3)}$	$\bar{v}_i$	2	$(\mathbf{3}, \mathbf{1}; \mathbf{1}, \mathbf{1})_{(1/6, -2/3)}$	$v_i$

# The location of Higgs bosons

Typically there could be a multitude of Higgs doublets (and triplets) in the spectrum

- triplets heavy or projected out
- exactly two Higgs doublet multiplets should remain light
- all other heavy

This is the so-called  $\mu$  problem

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The MiniLandscape identifies exactly one Higgs pair protected by a discrete R-symmetry and provides a unique solution to the  $\mu$  problem, because the

Higgs bosons live in the untwisted sector (delocalized Higgs as in torus compactification: remnants of  $N = 4$  Susy)

# Location of top quark

Given the fact that **the Higgs multiplets live in the bulk** we now explore how to obtain a large top quark Yukawa coupling

- need maximum “overlap” with the Higgs multiplet
- results of the MiniLandscape teach us that this requires the **top quark to live in the untwisted sector** as well



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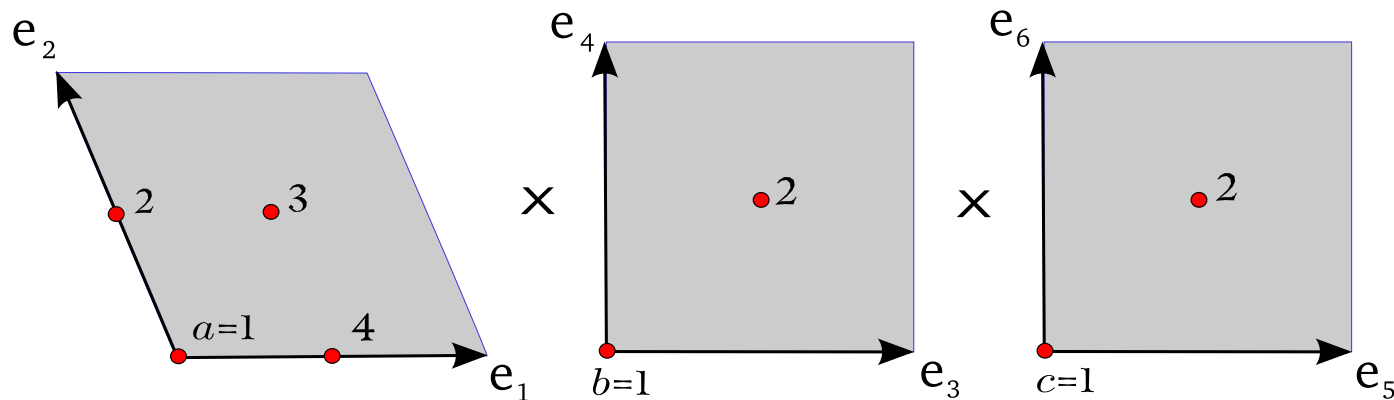
**Top quark in untwisted sector (bulk).** The third family is usually distributed over various sectors (it is not in a complete localized  $SO(10)$  representation).

Side remark:

**3 “complete” families impossible within  $Z_6II$  orbifold**

# First and second family

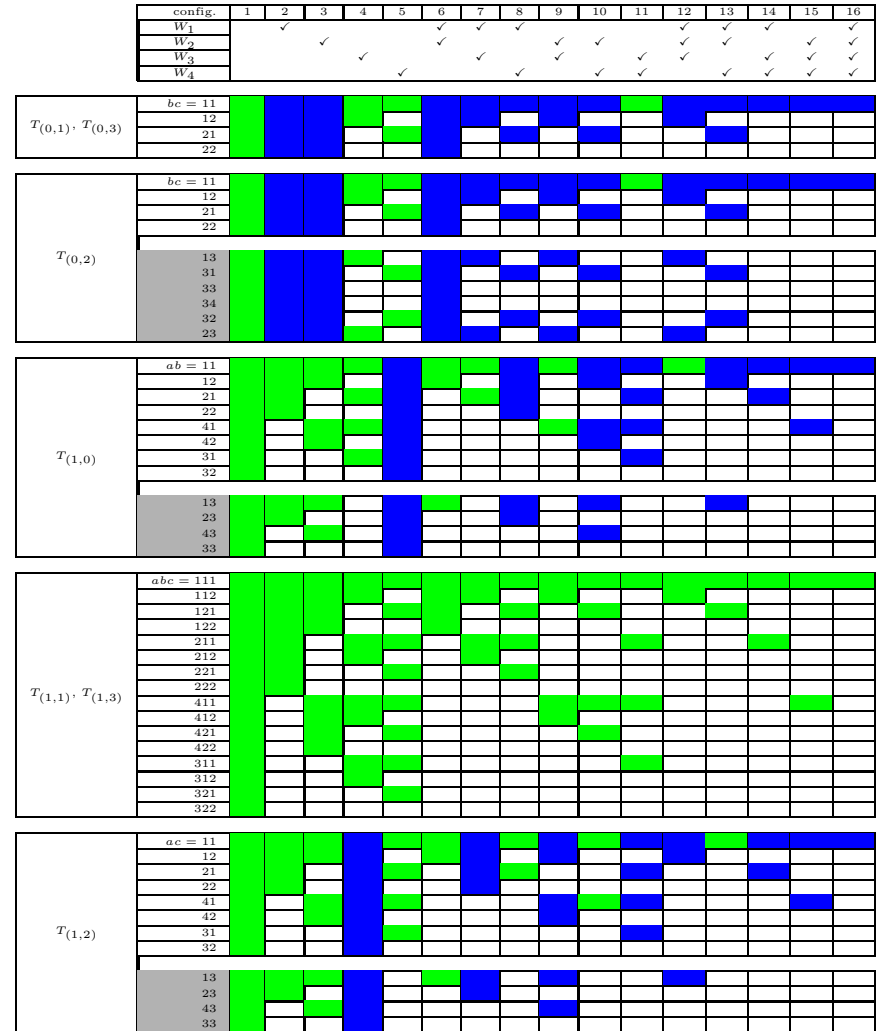
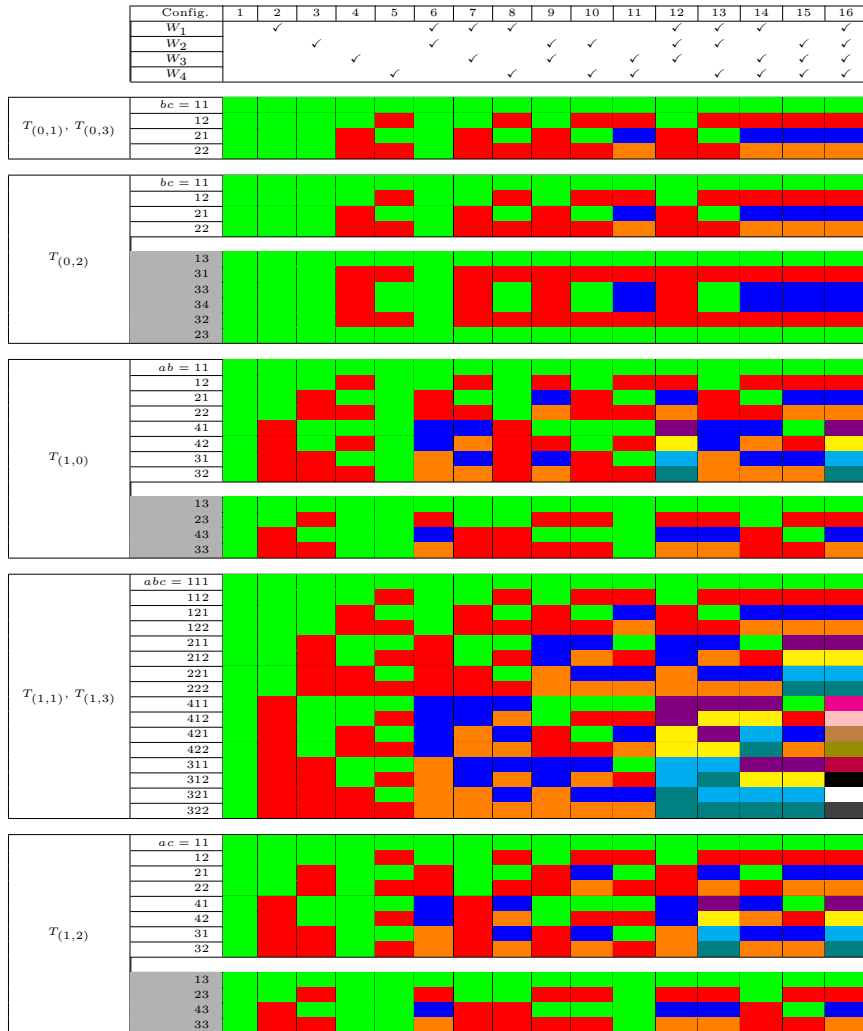
The first and second families are in complete localized 16-dimensional representation of  $SO(10)$  (at points of “enhanced” gauge symmetry)



They live in the  $\theta$  twisted sector and are localized at the fixed points  $a = 1$  and  $2$ ,  $b = 1$ ,  $c = 1$

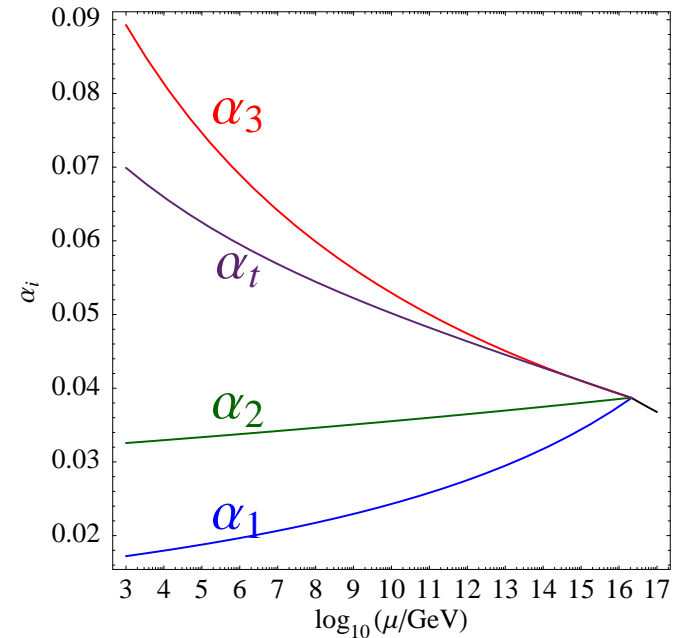
exhibiting a  $D_4$  family symmetry.

# Effects of Wilson lines



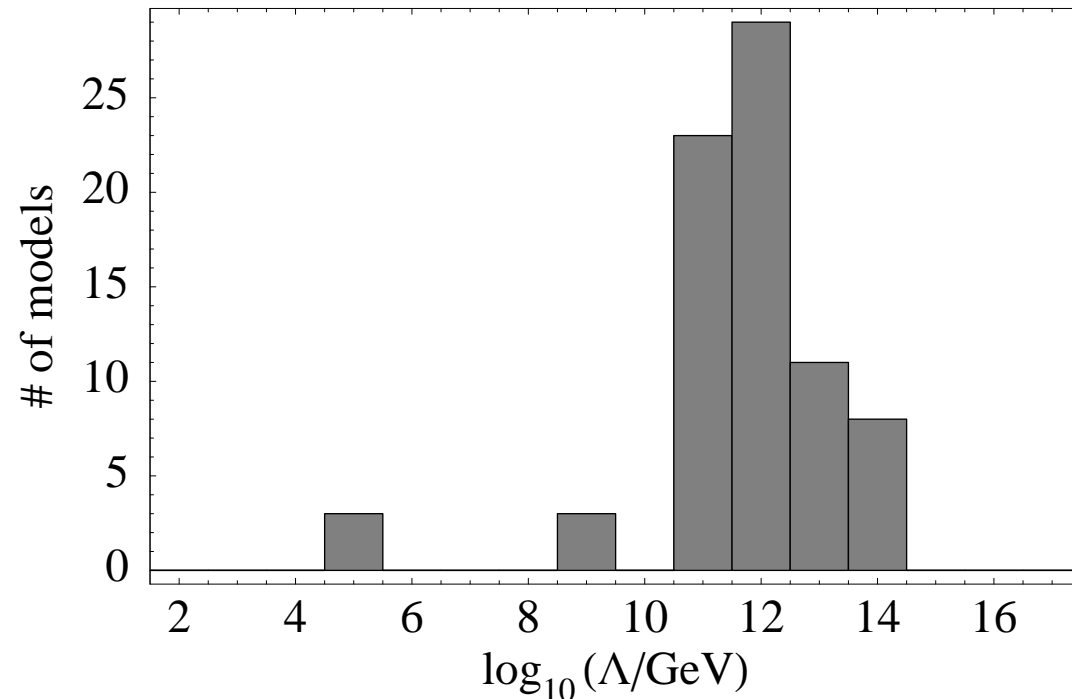
# Unification

- Higgs doublets live in the **bulk**
- **heavy top quark lives in the bulk as well.**
- $\mu$ -term protected by a **discrete R-symmetry**



- Minkowski vacuum before Susy breakdown (no AdS)
- **solution to  $\mu$ -problem** (Casas, Munoz, 1993)
- **first two families localized (smaller Yukawa couplings) exhibiting a discrete family symmetry**

# Heterotic string: gaugino condensation



Gravitino mass  $m_{3/2} = \Lambda^3 / M_{\text{Planck}}^2$  and  $\Lambda \sim \exp(-\tau)$

SU(4) in hidden sector predicts gravitino mass in TeV range

(Lebedev, HPN, Raby, Ramos-Sanchez, Ratz, Vaudrevange, Wingerter, 2006)

# Gaugino condensation 1985

CERN-TH.4123/85

ON THE LOW ENERGY  $d = 4$ ,  $N = 1$  SUPERGRAVITY THEORY

EXTRACTED FROM THE  $d = 10$ ,  $N = 1$  SUPERSTRING<sup>\*</sup>)

J.P. Derendinger, L.E. Ibáñez<sup>†)</sup>

CERN - Geneva

and

H.P. Nilles

Department of Theoretical Physics  
University of Geneva

# Gaugino condensation 1985

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## APPENDIX B

### KAHLER POTENTIAL AND DERIVATIVES IN THE TRUNCATED LAGRANGIAN

The Kähler potential upon truncation has the form

$$G \equiv K + \log |W|^2 = -\log(S+S^*) - 3\log(T+T^* - 2|c|^2) + \log |W|^2 \quad (\text{B.1})$$

From now on we shall define

$$s \equiv \text{Re} S, \quad t \equiv \text{Re} T, \quad t_c \equiv t - |c|^2 \quad (\text{B.2})$$

so that one has

$$K = -\log(2s) - 3\log(2t_c) \quad (\text{B.3})$$

# Gaugino condensation 1985

so that one has

$$K = -\log(2s) - 3\log(2t_c) \quad (\text{B.3})$$

Notice that the following useful identities hold

$$t_c = e^{\sigma} \varphi^{3/4} = g_4^2 e^{4\sigma} \quad (\text{B.4})$$

$$s = e^{3\sigma} \varphi^{-3/4} = g_4^{-2} \quad (\text{B.5})$$

$$\varphi = t_c s^{-1/3} \quad (\text{B.6})$$

$$e^{4\sigma} = s t_c \quad (\text{B.7})$$

$$e^K = \frac{1}{16} e^{-6\sigma} \varphi^{-3/2} = \frac{1}{16} s^{-1} t_c^{-3} \quad (\text{B.8})$$

The derivatives of the Kähler potential with respect to the scalar fields  $T$ ,  $S$ , and  $C$  are useful in order to obtain the explicit form of the  $N = 1$  supergravity Lagrangian. For the convenience of the reader we present many of these derivatives in this appendix. Denoting, for example,



# The overall pattern

The MiniLandscape provides a specific pattern for the soft masses with a large gravitino mass in the multi-TeV range

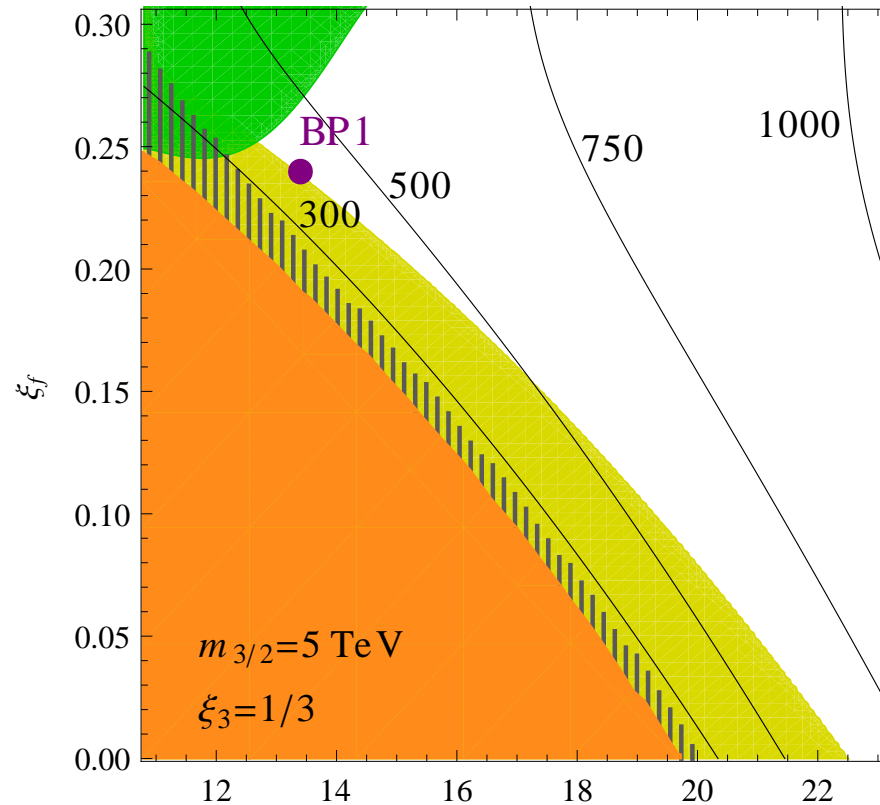
- normal squarks and sleptons in Multi-TeV range
- top squarks ( $\tilde{t}_L, \tilde{b}_L$ ) and  $\tilde{t}_R$  in TeV-range  
(suppressed by  $\log(M_{\text{Planck}}/m_{3/2}) \sim 4\pi^2$ )
- A-parameters in TeV range
- gaugino masses in TeV range
- mirage pattern for gaugino masses  
(compressed spectrum)
- heavy moduli (enhanced by  $\log(M_{\text{Planck}}/m_{3/2})$   
compared to the gravitino mass)

# A Closer Look

A more detailed picture requires the analysis of specific models. Issues that have to be clarified:

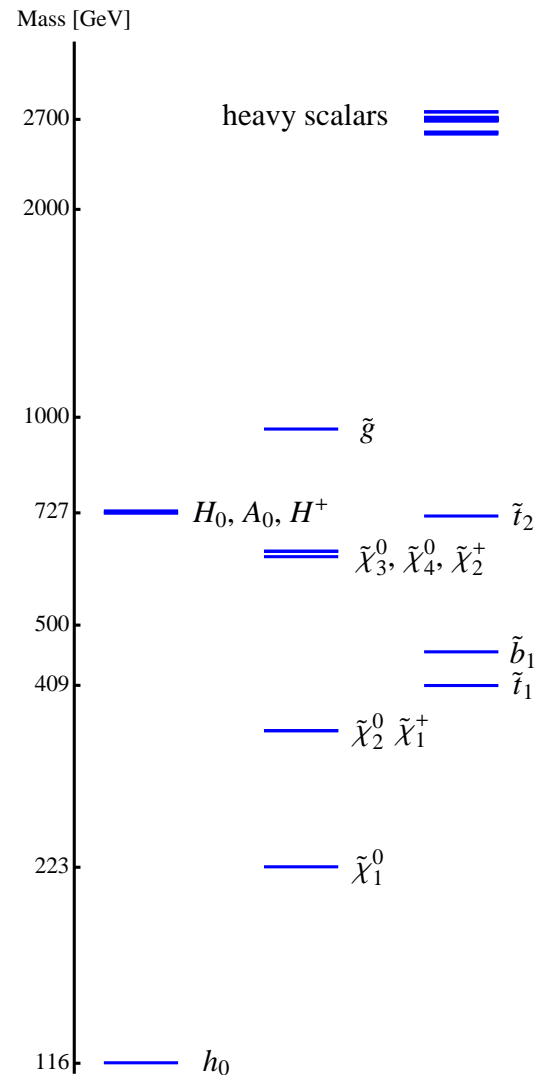
- the appearance of tachyons,
  - partially inherited from anomaly mediation
  - two loop effects in the presence of heavy scalars
- the hierarchy between gauginos and sfermions.
- Can we satisfy all phenomenological constraints?
  - mass of Higgs, correct electroweak symmetry breakdown etc.
  - nature and abundance of WIMP-LSP.
- What is the LHC reach to test this scheme?

# Benchmark model with a TeV gluino

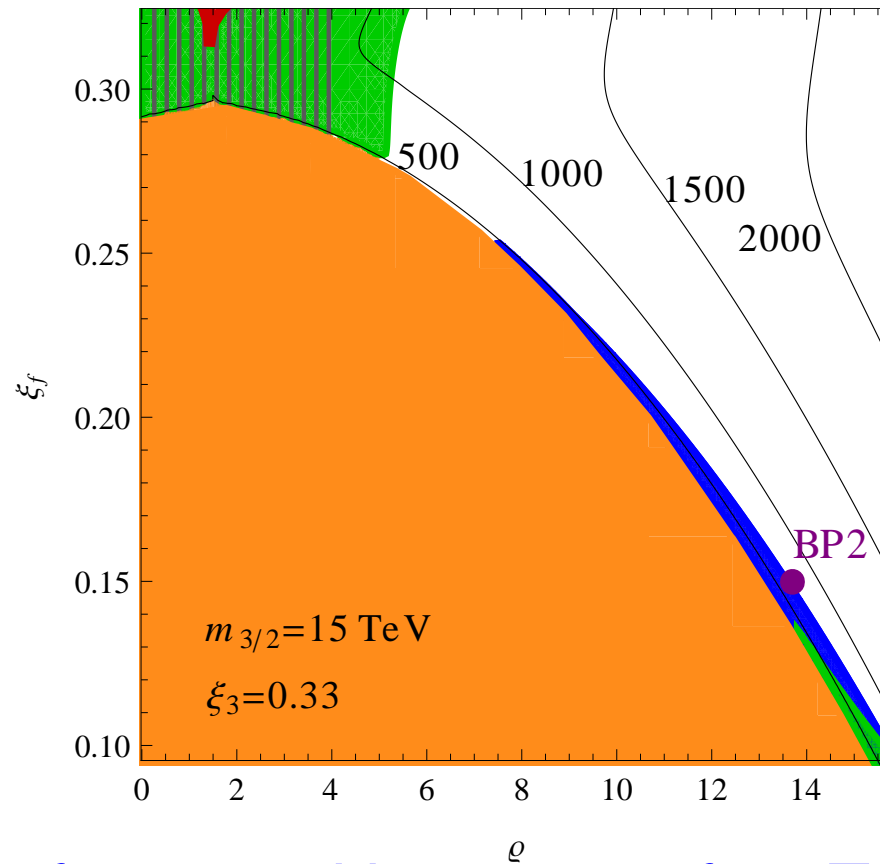


Parameter scan for a gluino mass of 1 TeV. The coloured regions are excluded while the hatched region indicates the current reach of the LHC. The contours indicate the mass of the lightest stop.

# Spectrum



# Model with a 3 TeV Gluino

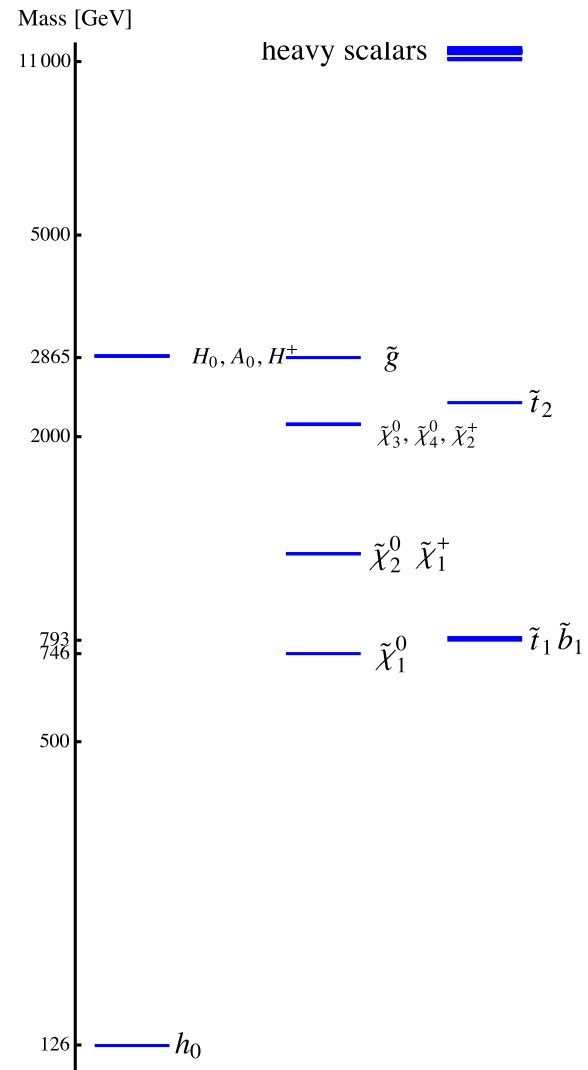


Parameter scan for a gravitino mass of 15 TeV.

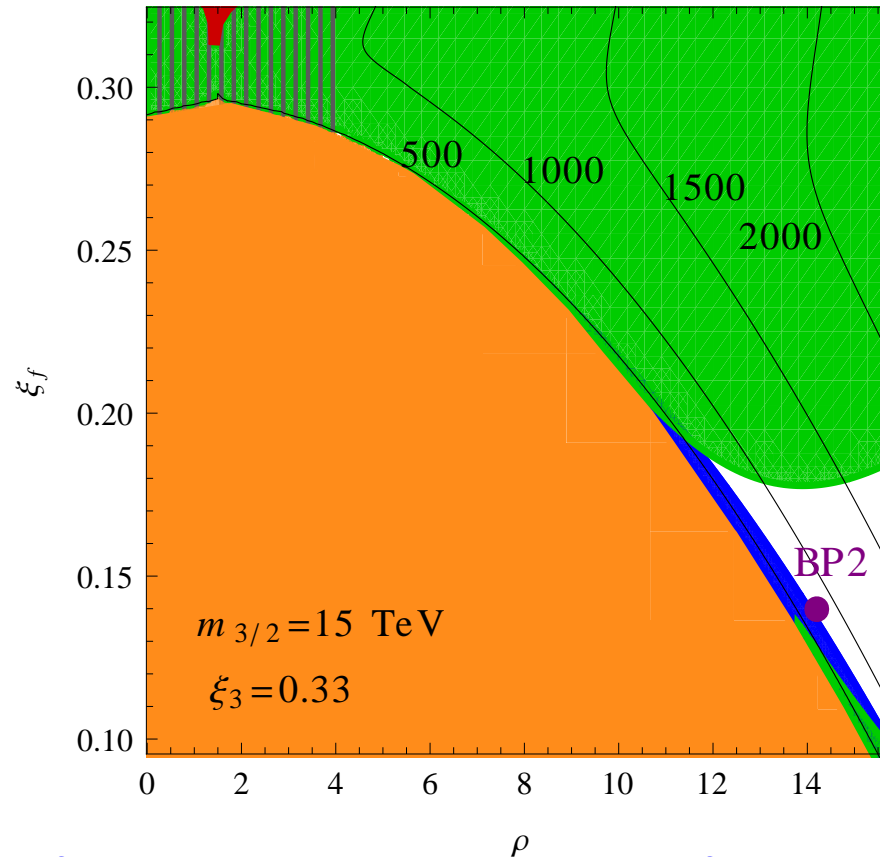
The coloured regions are excluded while the hatched region indicates the current reach of the LHC.

The contours indicate the mass of the lightest stop.

# Spectrum



# After Higgs discovery



Parameter scan for a gravitino mass of 15 TeV.

The coloured regions are excluded while the hatched region indicates the current reach of the LHC.

The contours indicate the mass of the lightest stop.

# Messages

- large gravitino mass (multi TeV-range)
- gaugino masses and stops suppressed by  $\log(M_{\text{Planck}}/m_{3/2})$
- other sfermion masses are of order  $m_{3/2}$
- the heterotic string yields “Natural SUSY” as a remnant of the underlying  $N = 4$  Susy
  - mirage pattern for gauginos,
  - light stop masses
- and this is a severe challenge for LHC searches.



# The overall scale

There is no (reliable) prediction for the gravitino mass

- except for fine-tuning arguments
- “no lose” criterion (SSC with 20+20 TeV)
- does LHC satisfy this criterion?

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Betting in the early 80's

- I bet that supersymmetry will be discovered before SSC gets into operation
- I bet that supersymmetry will have been forgotten before SSC gets into operation

# Conclusions

## Localization of quarks, leptons and Higgs bosons

- realistic models require Higgs multiplets and top multiplets in **bulk** (connected to  $\mu$  problem)
- this implies Gauge-Yukawa unification
- other fields tend to be localized at fixed points (tori) and exhibit discrete family symmetries

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## The legacy from extra dimensions ( $D = 10$ )

- discrete family symmetries
- mirage mediation (a hierarchy from  $\log(M_{\text{Planck}}/m_{3/2})$ )
- mass spectrum of **“Natural Susy”** from  $N = 4$

# Where does LHC lead us?



# Does he know?



# It seems as if.....?



# Mirage pictures

