
General Relativity

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1. Free movement on geodesics

- (a) Use the action principle $\delta \int \mathcal{L} dt = 0$ and conservation of energy to show that the movement of a free particle is always given by $\delta \int ds = 0$ (*i.e.* free particles move on a path of extremal length, the *geodesic*).
- (b) Write the equations of motion in last week's question 4 for a free particle on a sphere of fixed radius R .
- (c) Using the conserved conjugate momentum to ϕ , show that the remaining equation of motion can be written as

$$\ddot{\theta} - B^2 \frac{\cos \theta}{\sin^3 \theta} = 0, \quad (1)$$

where $B = \text{const} = \sin^2 \theta \dot{\phi}$. In the following, we will try to determine which geometrical form the motion will take.

- (d) Show that energy conservation implies

$$\frac{\dot{\theta}^2}{\sin^4 \theta \dot{\phi}^2} = \frac{A}{B^2} - \frac{1}{\sin^2 \theta}; \quad A = \text{const} = \frac{2E}{mR^2}. \quad (2)$$

- (e) Calculate $\frac{d}{d\phi} \cot \theta$ and substitute $u \equiv \cot \theta$ to rewrite (2) as

$$\left(\frac{du}{d\phi} \right)^2 = \frac{A}{B^2} - 1 - u^2. \quad (3)$$

- (f) Show that the solution of (3) can be rewritten in cartesian coordinates as

$$z = \alpha x + \beta y, \quad x^2 + y^2 + z^2 = R^2, \quad (4)$$

where α and β are suitably chosen constants.

What form do the trajectories of a free particle on a sphere take?

2. Riemann Tensor

The Christoffel symbols are *not* tensors, and thus are not suitable to describe a curved geometry in a coordinate-invariant way. The only tensor that can be constructed from the metric and its first and second derivatives is the *Riemann tensor*

$$R^\lambda{}_{\mu\nu\kappa} \equiv \frac{\partial \Gamma^\lambda{}_{\mu\nu}}{\partial x^\kappa} - \frac{\partial \Gamma^\lambda{}_{\mu\kappa}}{\partial x^\nu} + \Gamma^\eta{}_{\mu\nu} \Gamma^\lambda{}_{\kappa\eta} - \Gamma^\eta{}_{\mu\kappa} \Gamma^\lambda{}_{\nu\eta}. \quad (5)$$

Through self-contractions we get the *Ricci tensor* $R_{\mu\kappa} \equiv R^\lambda{}_{\mu\lambda\kappa}$ and the *curvature scalar* $R \equiv g^{\mu\kappa} R_{\mu\kappa}$.

(a) Using the metric, the Riemann tensor can be made fully covariant:

$$\begin{aligned} R_{\lambda\mu\nu\kappa} &= g_{\lambda\eta} R^\eta{}_{\mu\nu\kappa} \\ &= \frac{1}{2} \left(\frac{\partial^2 g_{\lambda\nu}}{\partial x^\kappa \partial x^\mu} - \frac{\partial^2 g_{\mu\nu}}{\partial x^\kappa \partial x^\lambda} - \frac{\partial^2 g_{\lambda\kappa}}{\partial x^\nu \partial x^\mu} + \frac{\partial^2 g_{\mu\kappa}}{\partial x^\nu \partial x^\lambda} \right) \\ &\quad + g_{\eta\sigma} (\Gamma^\eta{}_{\nu\lambda} \Gamma^\sigma{}_{\mu\kappa} - \Gamma^\eta{}_{\kappa\lambda} \Gamma^\sigma{}_{\mu\nu}). \end{aligned} \quad (6)$$

Check the symmetry properties $R_{\lambda\mu\nu\kappa} = -R_{\lambda\mu\kappa\nu}$ and $R_{\lambda\mu\nu\kappa} = +R_{\nu\kappa\lambda\mu}$.

- (b) Calculate the components of $R^\ell{}_{mnk}$, R_{mk} and the curvature scalar R for a space with coordinates (θ, ϕ) and metric $g_{mn} = \text{diag}(a^2, a^2 \sin^2 \theta)$.
(Use last week's Christoffel symbols again!)
- (c) Do the same for a space with coordinates (x^1, x^2) and metric $g_{mn} = \text{diag}(1, (x^1)^2)$.
What is the geometry?
- (d) What is the curvature of a circle with fixed radius R ?