
General Relativity

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1. The Einstein-Hilbert action

The field equations of General Relativity can be derived from a variational principle. Let's see how this works!

(a) Prove the four useful identities

$$\begin{aligned}(i) \quad \delta g^{\mu\nu} &= -g^{\mu\kappa} g^{\nu\lambda} \delta g_{\kappa\lambda}, \\(ii) \quad \delta g &= g g^{\mu\nu} \delta g_{\mu\nu}, \\(iii) \quad \delta \sqrt{-g} &= \frac{1}{2} \sqrt{-g} g^{\mu\nu} \delta g_{\mu\nu}, \\(iv) \quad \delta R_{\mu\nu} &= D_\nu \delta \Gamma^\kappa_{\kappa\mu} - D_\kappa \delta \Gamma^\kappa_{\mu\nu},\end{aligned}$$

where $g \equiv \det g_{\mu\nu}$. (*Hint: To prove (ii) use $\ln(\det g_{\mu\nu}) = \text{tr}(\ln g_{\mu\nu})$.*)

(b) Define the *Einstein-Hilbert* action by

$$S_{\text{EH}} \equiv \frac{1}{16\pi G} \int R \sqrt{-g} \, d^4x, \quad (1)$$

where the constant factor $\frac{1}{16\pi G}$ is introduced to reproduce the Newtonian limit when matter is added. Show that under the variation $g_{\mu\nu} = g_{\mu\nu} + \delta g_{\mu\nu}$

$$\delta S_{\text{EH}} = \frac{1}{16\pi G} \int \left(-R^{\mu\nu} + \frac{1}{2} R g^{\mu\nu} \right) \delta g_{\mu\nu} \sqrt{-g} \, d^4x.$$

(c) Suppose there exists matter described by an action

$$S_{\text{M}} \equiv \int \mathcal{L} \sqrt{-g} \, d^4x, \quad (2)$$

where \mathcal{L} is the Lagrangian density of the theory. If the matter action changes by δS_{M} under $\delta g_{\mu\nu}$, the *energy-momentum tensor* $T^{\mu\nu}$ is defined by

$$\delta S_{\text{M}} = \frac{1}{2} \int T^{\mu\nu} \delta g_{\mu\nu} \sqrt{-g} \, d^4x. \quad (3)$$

What is the expression for $T^{\mu\nu}$ in terms of \mathcal{L} ?

- (d) Write down the field equations from varying $S_{\text{EH}} + S_{\text{M}}$ with respect to the metric.
- (e) We may add an extra scalar to the scalar curvature without spoiling the invariance of the action. Add a constant Λ , called the *cosmological constant* and write down the modified field equations.

2. Schwarzschild solution (Episode I)

A general ansatz of writing a spherically symmetric and static line element is

$$ds^2 = B(r) dt^2 - A(r) dr^2 - r^2(d\theta^2 + \sin^2 \theta d\phi^2). \quad (4)$$

Traditionally, this has been called the *standard form*. We will now (and on the next exercise sheet) try to determine $A(r)$ and $B(r)$ from the Einstein equations.

- (a) What are the non-zero components of the metric tensor?
- (b) For all, who want to calculate some Christoffel symbols: Calculate all non-zero Christoffel symbols

$$\Gamma_{\mu\nu}^{\lambda} = \frac{1}{2}g^{\lambda\rho} \left(\frac{\partial g_{\rho\mu}}{\partial x^{\nu}} + \frac{\partial g_{\rho\nu}}{\partial x^{\mu}} - \frac{\partial g_{\mu\nu}}{\partial x^{\rho}} \right). \quad (5)$$

(You should find 8 distinct terms.)

- (c) And for all, who still want to calculate more: Calculate the terms of the Ricci tensor

$$R_{\mu\kappa} \equiv \frac{\partial \Gamma_{\mu\lambda}^{\lambda}}{\partial x^{\kappa}} - \frac{\partial \Gamma_{\mu\kappa}^{\lambda}}{\partial x^{\lambda}} + \Gamma_{\mu\lambda}^{\eta} \Gamma_{\kappa\eta}^{\lambda} - \Gamma_{\mu\kappa}^{\eta} \Gamma_{\lambda\eta}^{\lambda}. \quad (6)$$

(Hint: $R_{\mu\nu} = 0$ for $\mu \neq \nu$.)

...and we will continue these nasty calculations next week...