Abundance of Thermal WIMPs in Non-standard Cosmological Scenarios

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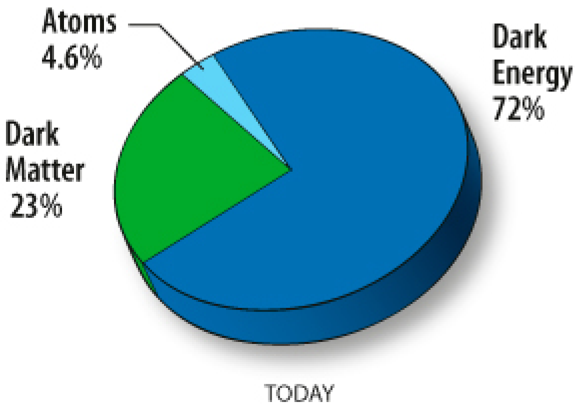
Refs:

- PRD73, 123502 (2006) [hep-ph/0603165]
- PRD76, 103524 (2007) [arXiv:0704.1590 [hep-ph]]

1. Motivation



[http://map.gsfc.nasa.gov]

- Observations of
  - cosmic microwave background
  - structure of the universe
  - etc.

Non-baryonic cold dark matter: $\Omega_{\rm DM}h^2 = 0.1099 \pm 0.0062$

- Weakly interacting massive particles (WIMPs) χ are good candidates for cold dark matter (CDM)

The predicted thermal relic abundance naturally explains the observed dark matter abundance: $\Omega_{\chi,{\rm standard}}h^2 \sim 0.1$

- Neutralino (LSP); 1st KK mode of the B boson (LKP); etc.

Investigation of early universe using CDM abundance

- The relic abundance of thermal WIMPs is determined by the Boltzmann equation: $\dot{n}_\chi + 3Hn_\chi = -\langle\sigma_{\rm eff}v\rangle(n_\chi^2 - n_{\chi,\rm eq}^2)$

  (and the reheat temperature: T_R)

- The (effective) cross section $\sigma_{\rm eff}$ can be (hopefully) determined from collider and DM detection experiments

→ We can test the standard CDM scenario and investigate conditions of very early universe: $T_R, H, \cdots$

- Standard scenario or non-standard scenarios?

  [Scherrer; Salati; Rosati; Profumo, Ullio; Pallis; Maisero, Pietroni, Rosati; ...]

Outline

This work

- We provide an approximate analytic treatment that is applicable to low-reheat-temperature scenarios
- Based on the assumption that dark matter is made of thermal WIMPs,
  - we derive the lower bound on the reheat temperature
  - we constrain possible modifications of the Hubble parameter

c.f. Cosmic p^+, γ ⇨ Bounds on pre-BBN expansion

[Schelke, Catena, Fornengo, Masiero, Pietroni PRD74 (2006); Donato, Fornengo, Schelke, JCAP0703 (2007)]

1. Motivation
2. Standard calculation of WIMP relic abundance
3. Low-temperature scenario
4. Constraints on the very early universe from WIMP dark matter
5. Summary

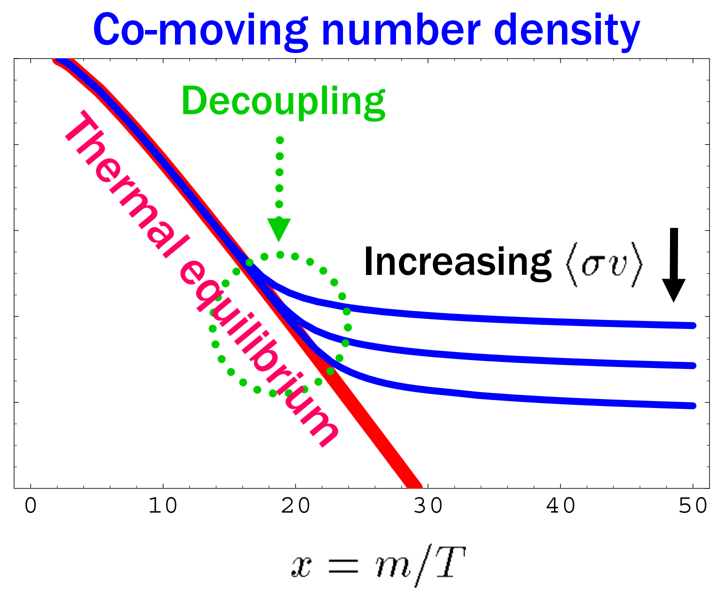
2. Standard calculation of the WIMP relic abundance

[Scherrer, Turner, PRD33(1986); Griest, Seckel, PRD43(1991); ...]

- Conventional assumptions:
  - $\chi = \bar{\chi}$, single production of χ is forbidden
  - Thermal equilibrium was maintained:
    $T_R(\text{Reheat temp}) \geq T_F(\text{Freezeout temp})$
- For adiabatic expansion the Boltzmann eq. is

$$\frac{\mathrm{d}Y_\chi}{\mathrm{d}x} = -\frac{\langle\sigma v\rangle s}{Hx}(Y_\chi^2 - Y_{\chi,\mathrm{eq}}^2),$$

$$Y_{\chi(,\mathrm{eq})} = \frac{n_{\chi(,\mathrm{eq})}}{s}, x = \frac{m_\chi}{T}$$



- χ decoupled when they were non-relativistic in RD epoch:

$$\langle\sigma v\rangle = a + 6b/x + \mathcal{O}(1/x^2), \quad n_{\chi,\mathrm{eq}} = g_\chi \ (m_\chi T/2\pi)^{3/2} \mathrm{e}^{-m_\chi/T}$$

$$\Omega_{\chi,\mathrm{standard}} h^2 \simeq 0.1 \times \left(\frac{a + 3b/x_F}{10^{-9}\ \mathrm{GeV}^{-2}}\right)^{-1} \left(\frac{x_F}{22}\right) \left(\frac{g_*}{90}\right)^{-1/2} \sim \Omega_{\mathrm{DM}} h^2$$

3. Low-temperature scenario

- **Assumptions:** $T_R \leq T_F; \; Y_\chi(x_0) = 0 \,, \; x_0 = m_\chi / T_R$
- **Zeroth order approximation:** For $T_R \ll T_F$, χ annihilation is negligible:

$$\frac{dY_0}{dx} = 0.028 \, g_\chi^2 g_*^{-3/2} m_\chi M_{\rm Pl} e^{-2x} x \left(a + \frac{6b}{x} \right)$$

The solution Y_0 is proportional to the cross section:

- **First order approximation:**
  - Add a correction term describing annihilation to Y_0 : $Y_1 = Y_0 + \delta \; (\delta < 0)$
  - As long as $|\delta| \ll Y_0$, the evolution equation for δ is

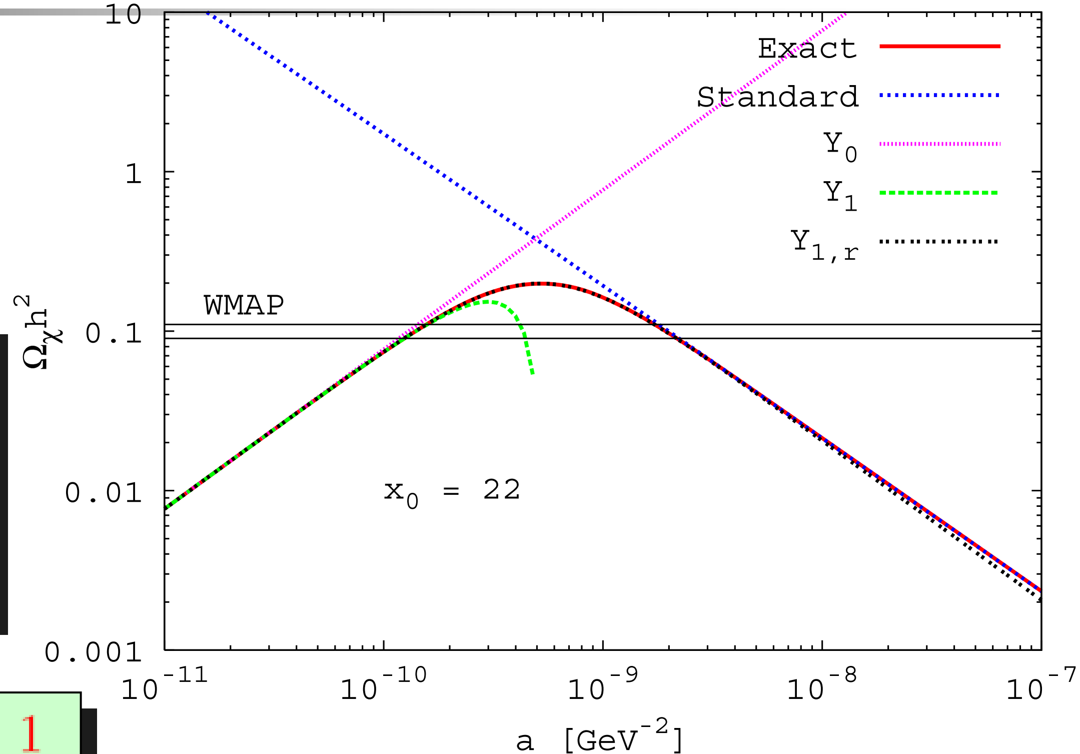
$$\frac{d\delta}{dx} = -1.3 \, \sqrt{g_*} m_\chi M_{\rm Pl} \left(a + \frac{6b}{x} \right) \frac{Y_0(x)^2}{x^2}$$

An analytic expression for δ can be obtained and is proportional to σ^3

Resummed ansatz

- We can obtain an excellent analytic formula for $T_R \leq T_F$ through resummation:

$$Y(x) = Y_0 + \delta = Y_0\left(1 + \frac{\delta}{Y_0}\right) \simeq \frac{Y_0}{1-\delta/Y_0} \equiv Y_{1,r}$$



- For $|\delta| \gg Y_0$, $Y_{1,r}(x) \simeq -\frac{Y_0^2}{\delta} \propto \frac{1}{\sigma}$

$x_0 \to x_F$ Standard formula

At late times, $Y_{1,r}(x \to \infty) = \frac{x_0}{1.3\sqrt{g_*}m_\chi M_{\rm Pl}(a + 3b/x_0)}$

4. Constraints on the very early universe from WIMP DM

- Out-of-equilibrium case: $\sigma \nearrow \Rightarrow \Omega h^2 \nearrow$; $T_R = m_\chi / x_R \nearrow \Rightarrow \Omega h^2 \nearrow$

  Equilibrium case: $\sigma \nearrow \Rightarrow \Omega h^2 \searrow$; $\Omega_\chi h^2$ is independent of T_R

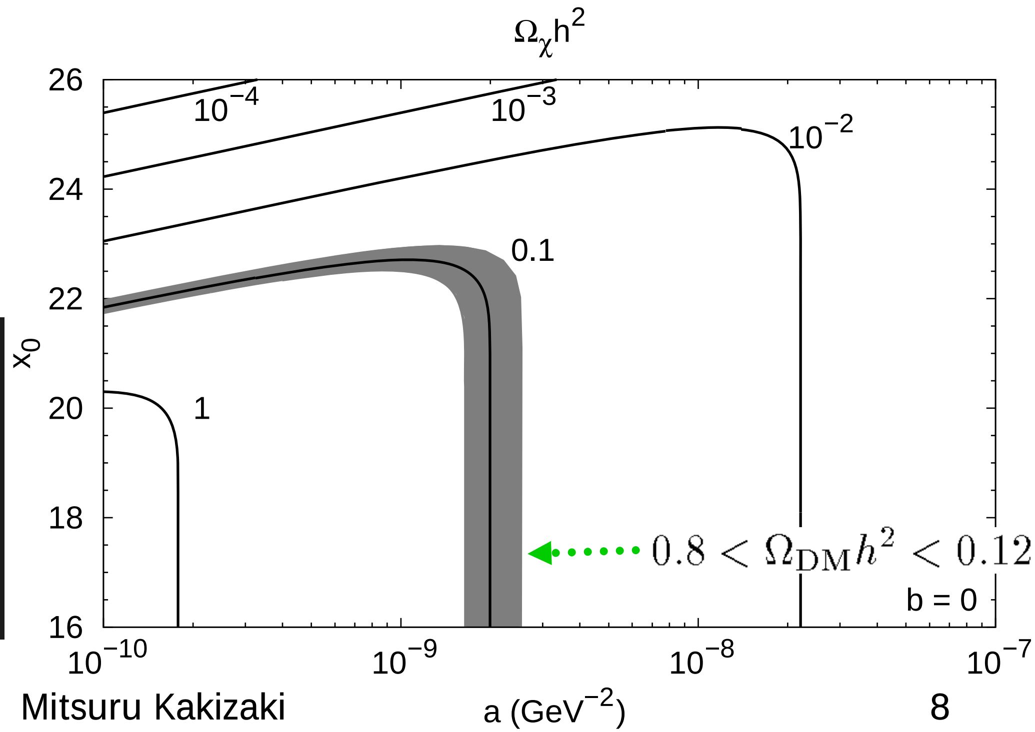
- Thermal relic abundance:

Assumption:
WIMPs are assumed to be produced in the RD epoch

Requirement that $\Omega_\chi h^2 \simeq 0.1$

↓

Lower bound on the reheat temperature: $T_R > m_\chi / 23$



Modified expansion rate

- Various cosmological models predict a non-standard early expansion

  ⇨ Predicted WIMP relic abundances are also changed

- Idea:

  Once we know $\sigma, \Omega_{\rm DM}h^2$

  ⇨ Constraint on the expansion rate at around WIMP decoupling (within the framework of thermal WIMP production)

- Parametrization: $A(z) = H_{\rm st}(z)/H(z)\,, \quad z \equiv T/m_\chi = 1/x$

- We need to know $A(z)$ only for $z_{\rm BBN} = 10^{-5} - 10^{-4} \le z \le z_F \sim 1/20 \ll \mathcal{O}(1)$

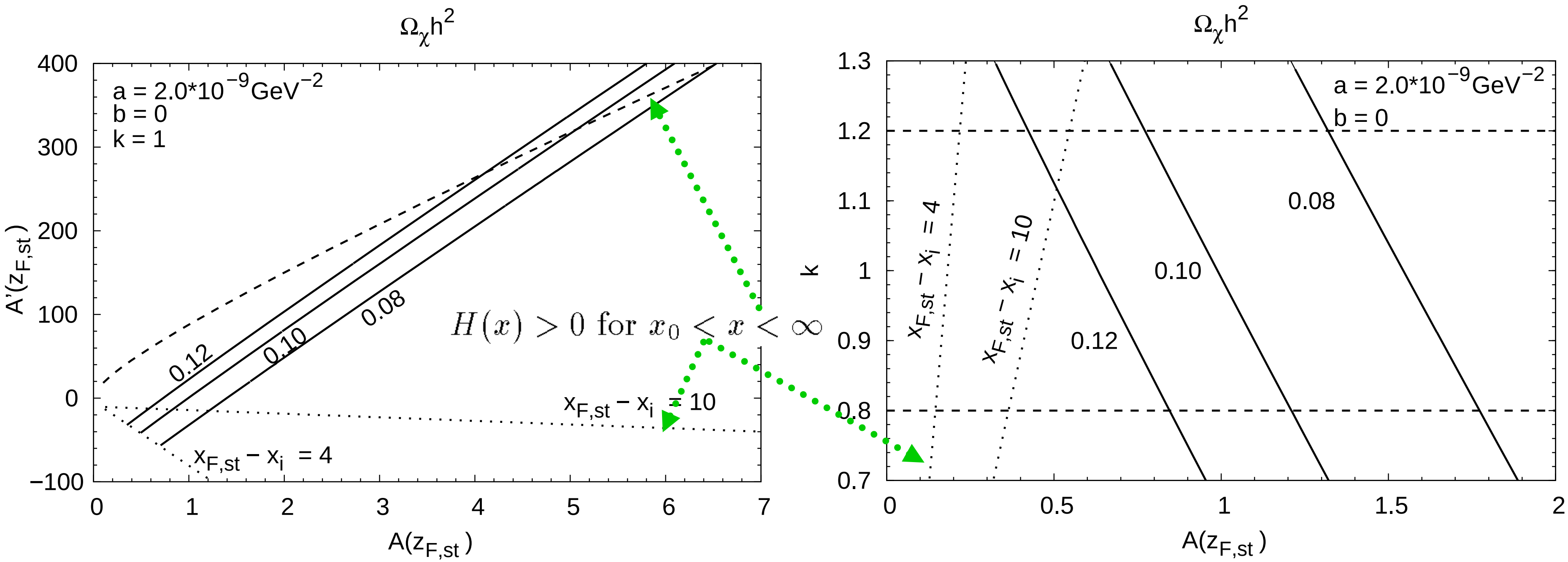
  ⇨ Taylor expansion of $A(z)$ in powers of $(z - z_{F,{\rm st}})$:

$$A(z) = A(z_{F,{\rm st}}) + (z - z_{F,{\rm st}})A'(z_{F,{\rm st}}) + \frac{1}{2}(z - z_{F,{\rm st}})^2 A''(z_{F,{\rm st}})$$

Subject to the BBN limit: $k \equiv A(z \to z_{\rm BBN}) \simeq 1$

Constraints on modifications of the Hubble parameter

- **Case for** $k \equiv A(z \to z_{\rm BBN}) = 1$
- **Case for** $A''(z_{F,\rm st}) = 0$



x_i : **Maximal temperature below which our expansion is valid**

$\Omega_\chi h^2$ **depends on all** $H(T_{\rm BBN} < T < T_F)$ ➡ **Weaker constraints on** $H(T_F)$

5. Summary

- Using the CDM relic density we can examine very early universe at around $T \sim m_\chi/20 \sim \mathcal{O}(10)\ \mathrm{GeV}$ (well before BBN: $T_{\mathrm{BBN}} \sim \mathcal{O}(1)\ \mathrm{MeV}$)

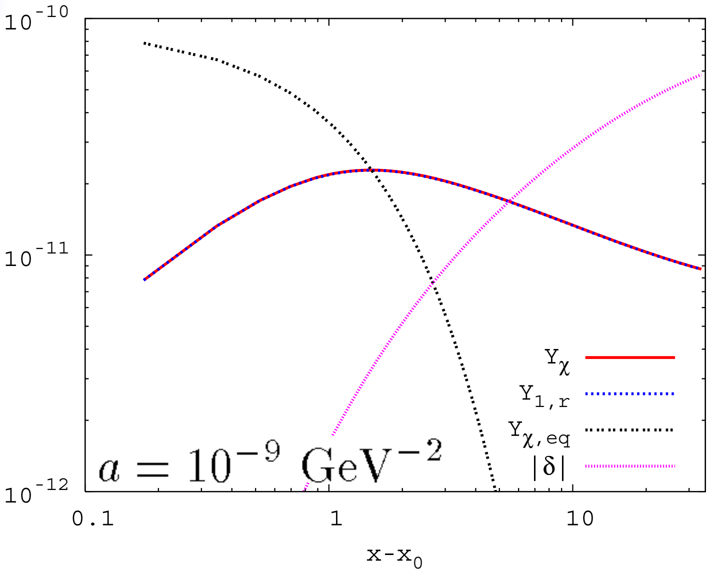
- We find an approximate analytic formula that is valid for all $T_R \leq T_F$

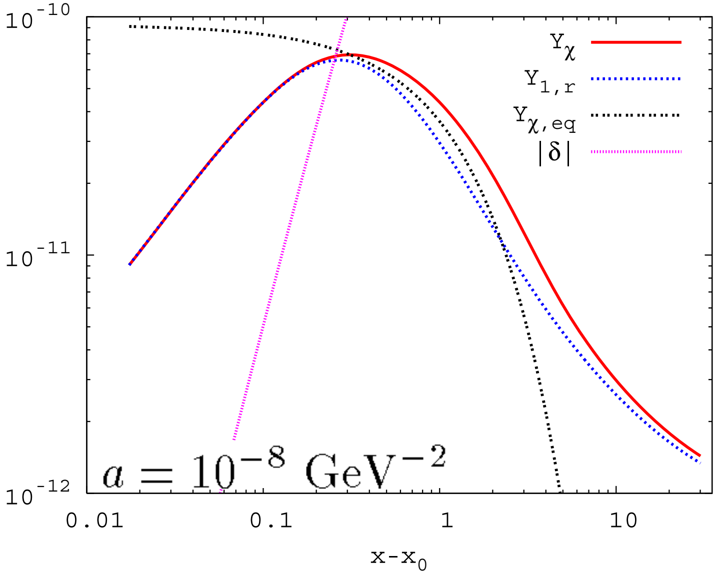
- By requiring $\Omega_{\chi,\mathrm{thermal}} h^2 = \Omega_{\mathrm{DM}} h^2$,
  we found the lower bound on the reheat temperature: $T_R > m_\chi/23$

- The sensitivity of $\Omega_{\chi,\mathrm{thermal}} h^2$ on $H(T_F)$ is weak
  because $\Omega_{\chi,\mathrm{thermal}} h^2$ depends on all $H(T_{\mathrm{BBN}} < T < T_F)$

Backup slides

Evolution of solutions

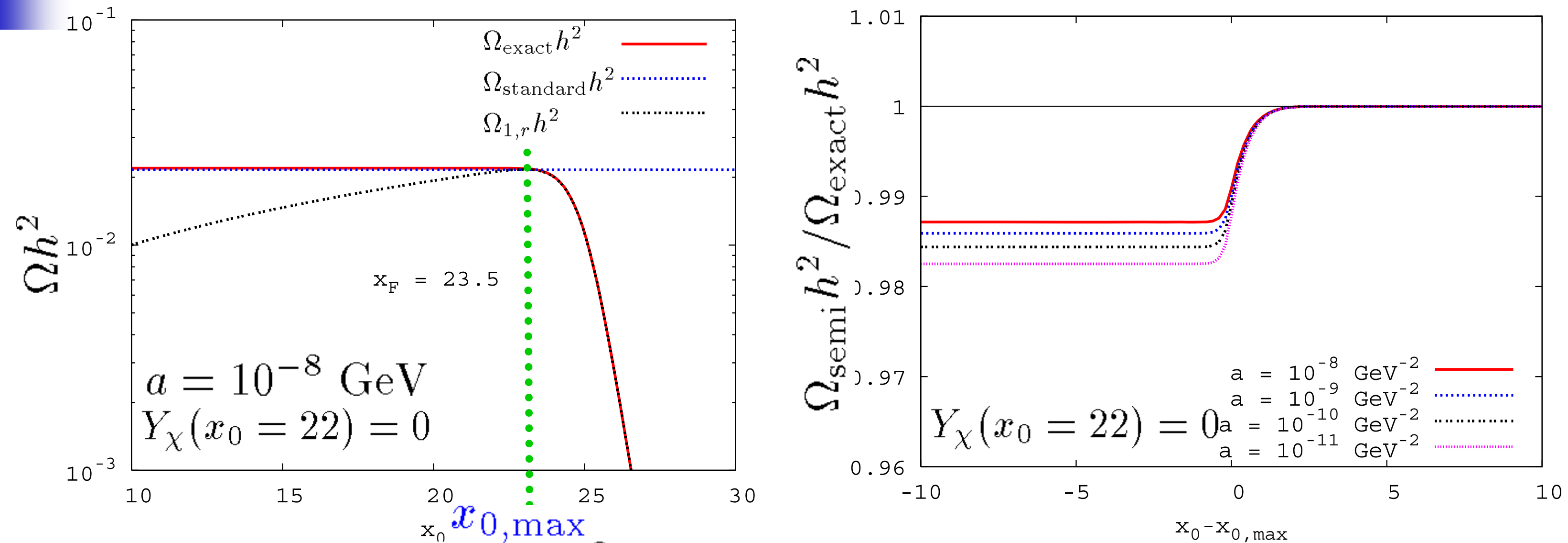




Y_χ : Exact result, $Y_{1,r}$: Re-summed ansatz, $b = 0,\ Y_\chi(x_0 = 22) = 0$

- The re-summed ansatz $Y_{1,r}$ describes the full temperature dependence of the abundance when equilibrium is not reached
- For larger cross section the deviation becomes sizable for $x - x_0 \sim 1$, but the deviation becomes smaller for $x \gg x_0$

Semi-analytic solution



- $Y_{1,r}(x_0, x \to \infty)$ $(\propto \Omega_{1,r} h^2)$ has a maximum (left)
- New semi-analytic solution can be constructed: $\Omega_{\rm semi} h^2$ (right)

For $x_0 > x_{0,\rm max}$, use $Y_{1,r}(x_0)$; for $x_0 < x_{0,\rm max}$, use $Y_{1,r}(x_{0,\rm max})$

The semi-analytic solution $\Omega_{\rm semi} h^2$ reproduces the correct final relic density $\Omega_{\rm exact} h^2$ to an accuracy of a few percent