

Collider signals and neutralino dark matter detection in relic-density-consistent models without soft term universality

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: [JHEP05 \(2008\) 058](#)

PLANCK 2008 - From the Planck Scale to the ElectroWeak Scale

May 19-23, 2008

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 - ★ Non-universal scalar mass models
 - ★ Non-universal gaugino mass models
- Implications for collider searches
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- Conclusions

Neutralino Dark Matter

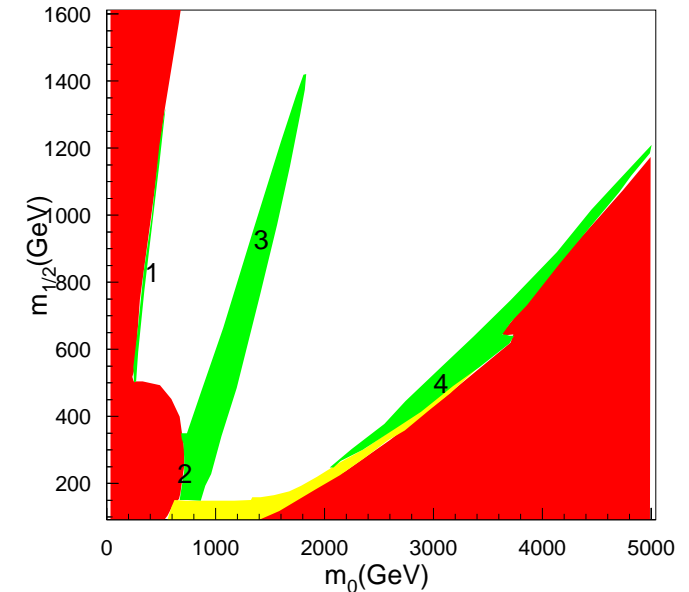
- Dominant composition of matter in our universe is not detected visibly but inferred from gravitational effects (Galactic Clustering, Rotation Curves, Gravitational Lensing, Cosmic Microwave Background ...)
- Dark Matter should be non-baryonic (no candidate in the SM), non-relativistic (cold), stable(or long-lived), weakly (or super-weakly) interacting matter
- In SUSY models with R -parity conservation
 - ⇒ the Lightest Supersymmetric Particle(LSP) is absolutely stable
 - ⇒ lightest neutralino \tilde{Z}_1 is the LSP in most of MSSM parameter space
 - ⇒ \tilde{Z}_1 is good candidate for Cold Dark Matter (CDM)
- From the WMAP results, the cold dark matter density of the universe is $\Omega_{CDM}h^2 = 0.111_{-0.015}^{+0.011}$: (upper bound is a tight constraint on SUSY models containing DM candidates : DM may consist of several components)
- Number density is governed by Boltzmann equation,

$$dn/dt = -3Hn - \langle\sigma v_{rel}\rangle(n^2 - n_0^2)$$
 - ⇒ requires evaluating many thousands Feynman diagrams
 - ⇒ high (co-)annihilation cross section implies low relic abundance

Motivations for SUSY models with non-universal soft terms

- **WMAP allowed Regions** in parameter space of mSUGRA

1. $\tilde{\tau}$ co-annihilation region at low m_0 , $m_{\tilde{\tau}_1} \sim m_{\tilde{Z}_1}$
2. bulk region at low m_0 and $m_{1/2}$, light sleptons (LEP2 excluded)
3. Higgs-funnel H, A resonance ($2m_{\tilde{Z}_1} \simeq m_{A,H}$) at large $\tan\beta \sim 50$ or h -resonance at low $m_{1/2}$ ($2m_{\tilde{Z}_1} \simeq m_h$)
4. FP/HB region at large m_0 , low $\mu \rightarrow$ mixed higgsino dark matter (MHDM)



- **Motivations for models with non-universality**

- ★ all relic-density-consistent regions in mSUGRA are near the edges of **theoretically**(or **experimentally**) excluded
- ★ need to examine how already drawn conclusions from the mSUGRA model are affected by relaxing the universality assumptions
- ★ within R -parity conserved neutralino dark matter assumption, WMAP value provides a strong constraint reducing model parameter space by one unit

Assumptions and Tools

- **Assumptions**

- MSSM is an effective theory between the weak and GUT scale
- R -parity is conserved
- Neutralino LSP
- (near)degeneracy of first and second generation of SSB sfermions \rightarrow FCNC suppressed
- CP-violating phases in SSB suppressed \rightarrow CP contribution of SUSY is small

- **Tools**

- all mass spectrum: ISAJET 7.76
- relic density, direct detection rate: IsaTools package (IsaReD, IsaReS)
- all indirect dark matter detection rates: DarkSUSY

- **Relic-density-consistent models** obtained by adjusting

- composition of neutralino (**WTN**: Well-Tempered Neutralino^{*})

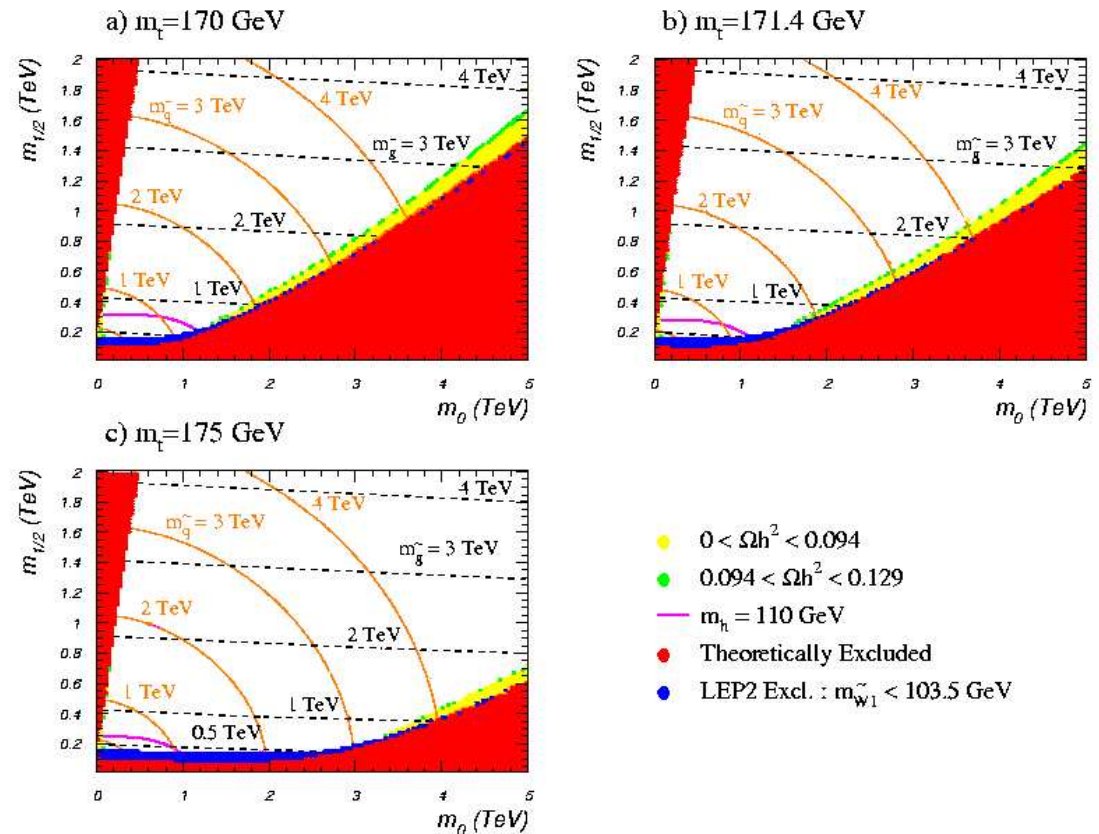
^{*}termed by Arkani-Hamed et al. Nucl.Phys.B741, 108, 2006

- masses of neutralino or other sparticles

Updated results on parameter space of mSUGRA 1

- **parameter Space**
 $m_0, m_{1/2}, A_0, \tan\beta, \text{sign}(\mu)$
- **effect of m_t variation**
 - location of EWSB excluded region
 → HB/FP region moves
 - gluino and squark mass contours hardly change
 - as m_t increased, A -funnel moves to smaller m_0

mSUGRA : $\tan\beta=10, A_0=0, \mu>0$



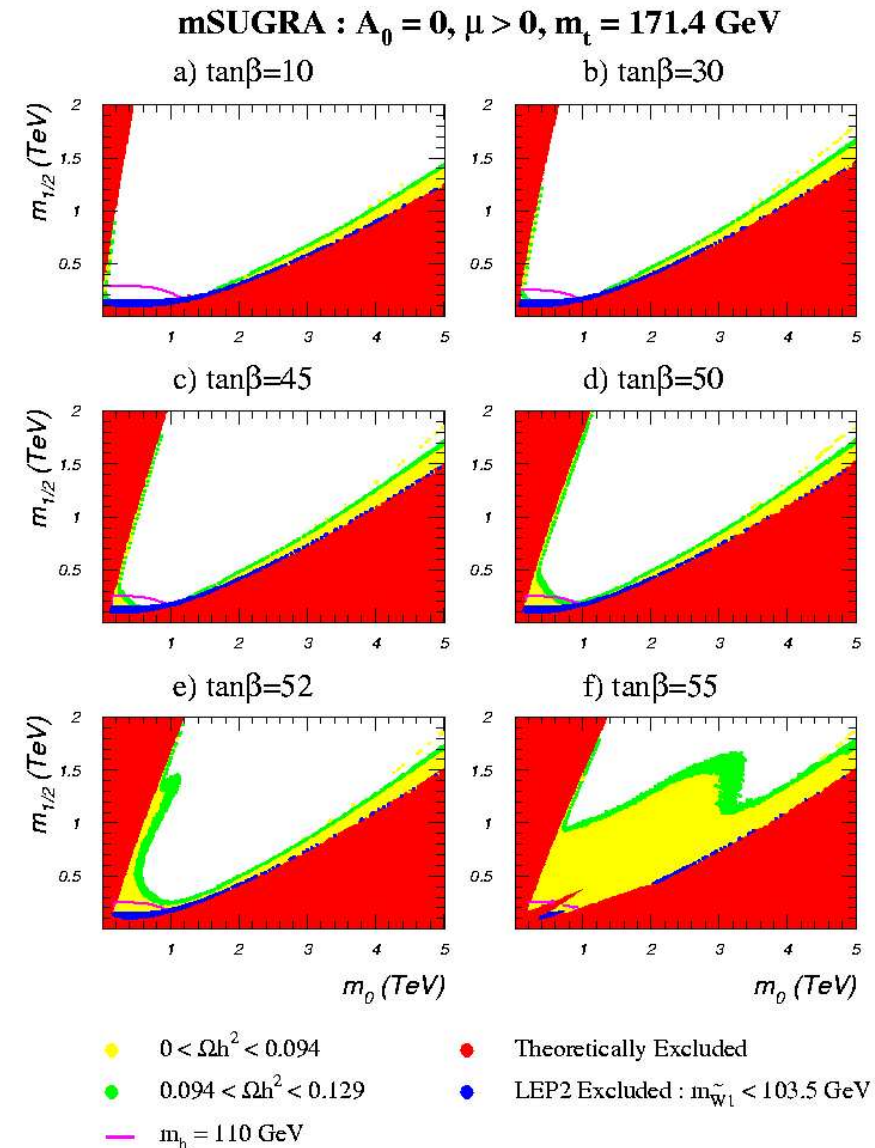
Updated results on parameter space of mSUGRA 2

- effect of $\tan\beta$ variation

- larger $\tan\beta$, wider WMAP allowed region
(as $\tan\beta$ grows, m_A drops due to b and τ Yukawa coupling effect)
- as $\tan\beta$ increased, A -funnel moves to larger m_0

- effect of $\mu < 0$

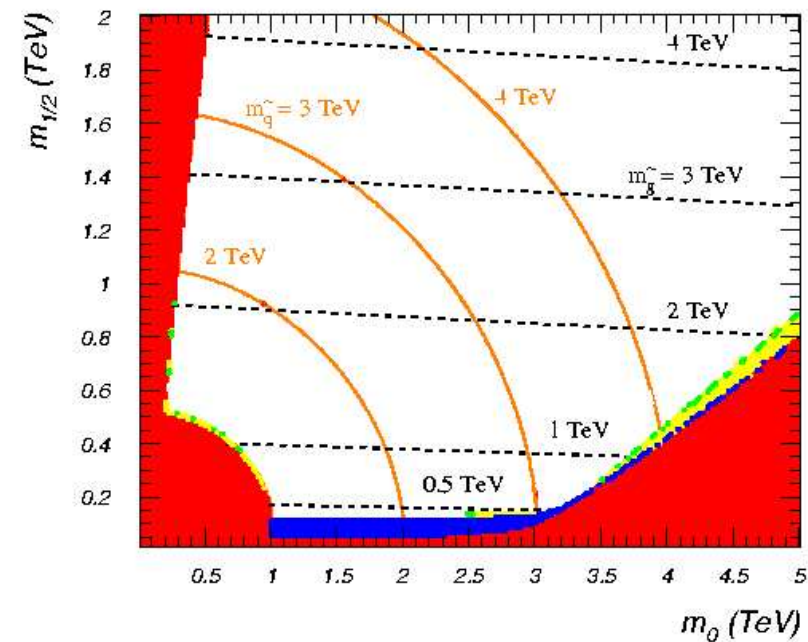
- A -funnel arises at lower $\tan\beta$
- narrower A -funnel



Updated results on parameter space of mSUGRA 3

- effect of large negative A_0
 - reduce $m_{\tilde{t}_1}$
 - $\tilde{t}_1 - \tilde{Z}_1$ mass gap quite small
 - relic density reduced via stop co-annihilation

mSUGRA : $\tan\beta=10, A_0=-2 \text{ TeV}, \mu>0, m_t=171.4 \text{ GeV}$



- $0 < \Omega h^2 < 0.094$
- $0.094 < \Omega h^2 < 0.129$
- Theoretically Excluded
- LEP2 Excl. : $m_{\tilde{W}_1} < 103.5 \text{ GeV}$

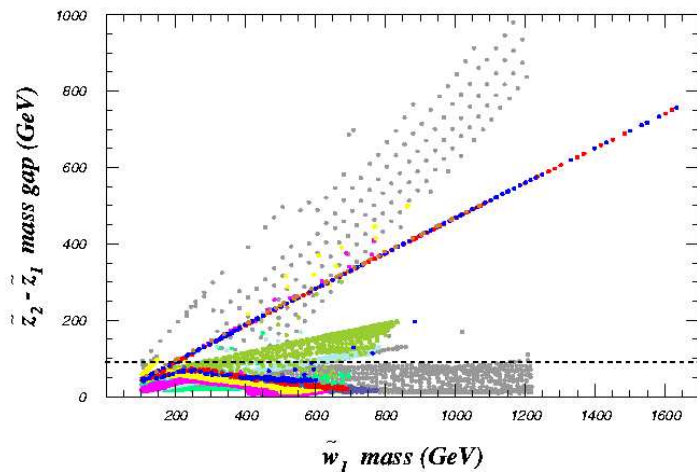
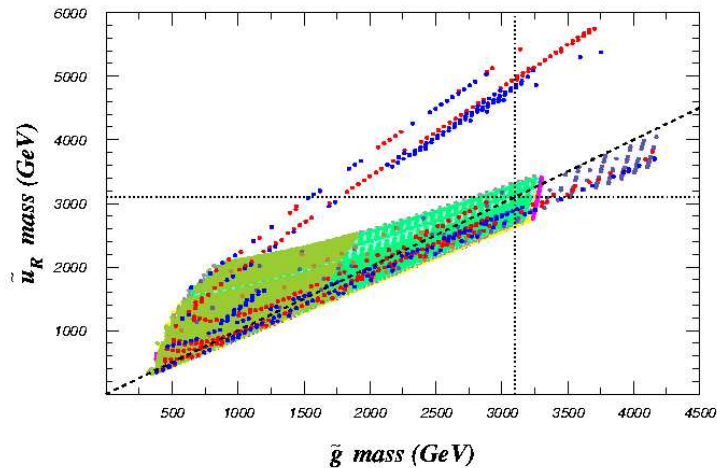
Non-universal scalar mass models

- generation non-universality: Normal scalar Mass Hierarchy (NMH)
 $m_0(1, 2), m_0, m_{1/2}, A_0, \tan\beta, \text{sign}(\mu)$
 - $m_0(1, 2)$: first/second generation, $m_0(3) = m_{H_u} = m_{H_d} \equiv m_0$: remaining
 - dial $m_0(1, 2)$ to low enough to bulk (co-)annihilation via light sleptons
- non-universal Higgs mass: one extra parameter case (NUHM1 $_{\mu}$, NUHM1 $_A$)
 $m_0, \delta_\phi, m_{1/2}, A_0, \tan\beta, \text{sign}(\mu)$
 - $m_\phi = m_0(1 + \delta_\phi), m_{H_u}^2 = m_{H_d}^2 \equiv \text{sign}(m_\phi)|m_\phi|^2$
 - $m_\phi > m_0$: small μ and MHDM
 - $m_\phi < 0$: $m_A \sim 2m_{\tilde{Z}_1} \rightarrow$ at any $\tan\beta$
- non-universal Higgs mass: two extra parameter case (HS-Higgs Splitting)
 $m_0, m_{H_u}^2$ (equivalently μ), $m_{H_d}^2$ (equivalently m_A), $m_{1/2}, A_0, \tan\beta, \text{sign}(\mu)$
 - $m_{H_{u,d}}^2 = m_0^2 (1 \mp \delta_H)$
 - $\delta_H < 0$: low μ and low m_A
 - $\delta_H > 0$: WMAP region via $\tilde{l}_L/\tilde{\nu}$ or \tilde{u}_R/\tilde{c}_R co-annihilation

Non-universal gaugino mass models

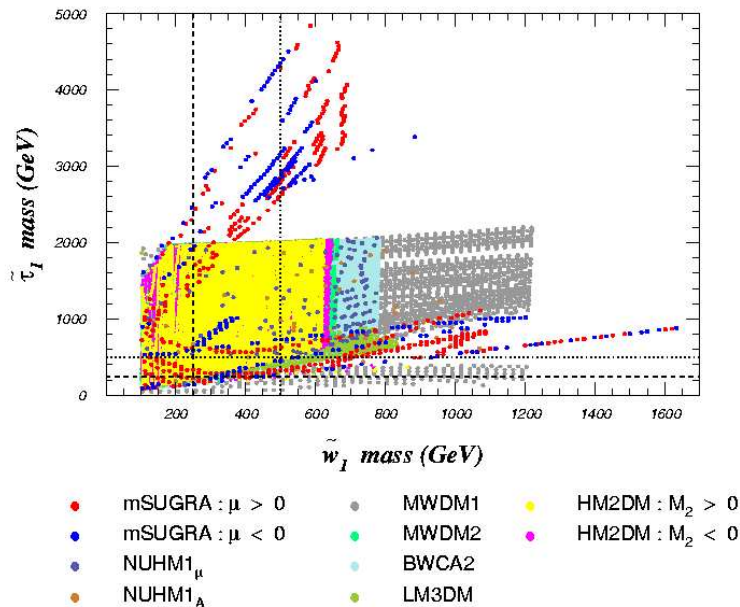
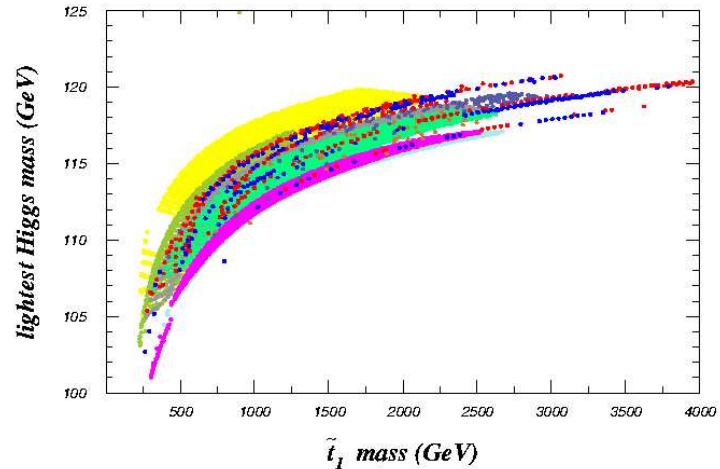
- Mixed Wino Dark Matter (**MWDM1**, **MWDM2**):
 $m_0, M_1(\text{or } M_2), m_{1/2}, A_0, \tan\beta, \text{sign}(\mu)$
 - by increasing the wino content of the LSP by reducing the ratio M_2/M_1
 - $M_1 \neq M_2 = M_3 = m_{1/2}$ or $M_2 \neq M_1 = M_3 = m_{1/2}$
- Bino-Wino Co-Annihilation Scenario (BWCA1, BWCA2):
 same as MWDM but M_1 and M_2 are in opposite sign
 - by allowing co-annihilation between high bino-like and wino-like states
- Low $|M_3|$ Dark Matter: Compressed SUSY (**LM3DM**):
 $m_0, M_3, m_{1/2}, A_0, \tan\beta, \text{sign}(\mu)$
 - by increasing the higgsino content of the LSP by decreasing the gluino mass
 - $M_3 \neq M_1 = M_2 = m_{1/2}$
- High $|M_2|$ Dark Matter: left-right split SUSY (**HM2DM**):
 $m_0, M_2, m_{1/2}, A_0, \tan\beta, \text{sign}(\mu)$
 - by allowing large M_2 mass
 - $M_2 \gg M_1 = M_3 = m_{1/2}$

Implications for collider searches 1



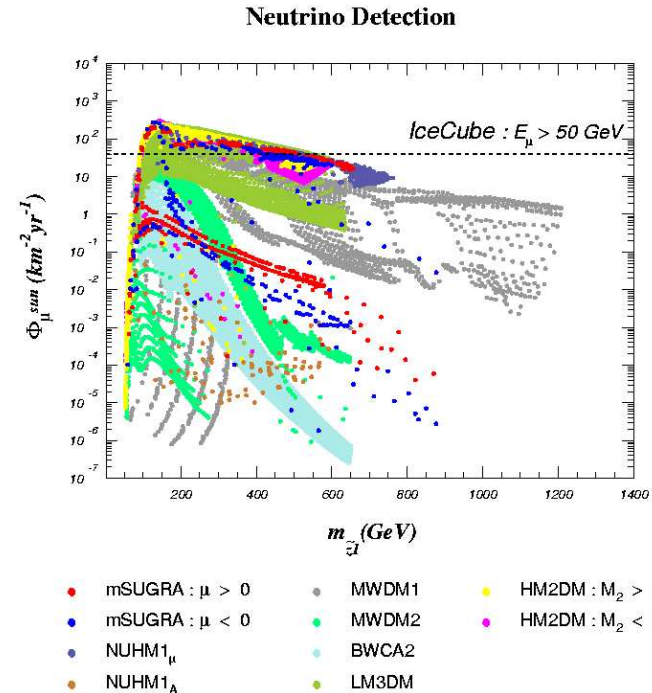
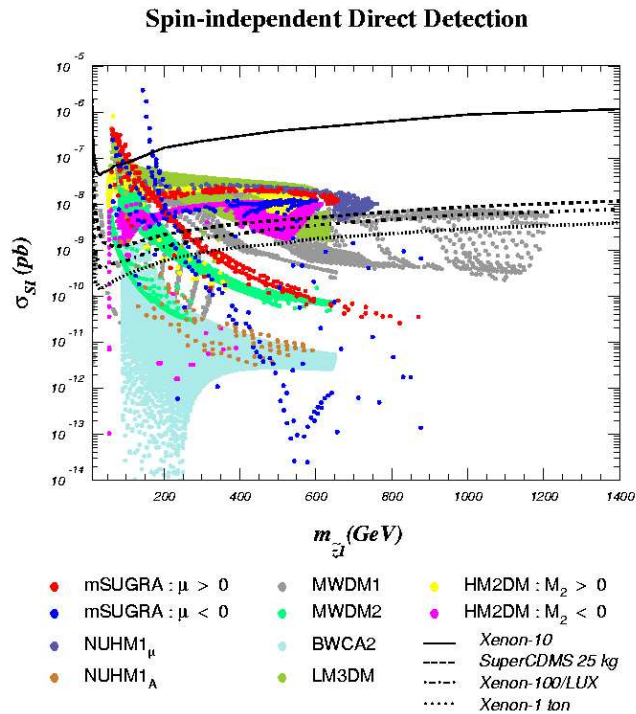
- with $A_0 = 0$, $m_t = 171.4$ GeV, $\tan \beta = 10$
 (except for the mSUGRA model: $\tan \beta = 10, 30, 45, 50, 52$ and 55)
- non-universal mass dialed to yield $\Omega_{\tilde{Z}_1} h^2 \simeq 0.11$
- $m_{\tilde{g}}$ vs. $m_{\tilde{u}_R}$
 - dotted lines: 100 fb^{-1} reach of CERN LHC
 - dashed line: $m_{\tilde{u}_R} = m_{\tilde{g}}$
 - most of models within reach of LHC except HB/FP region of mSUGRA
- $m_{\tilde{W}_1}$ vs. $m_{\tilde{Z}_2} - m_{\tilde{Z}_1}$
 - dashed line: $m_{\tilde{Z}_2} - m_{\tilde{Z}_1} = M_Z$
 - below the line, 3-body decay like $\tilde{Z}_2 \rightarrow \tilde{Z}_1 l \bar{l}$ open
 - in most models, $m(l\bar{l})$ mass edge visible at LHC

Implications for collider searches 2



- m_h vs. $m_{\tilde{t}_1}$
 - heavier \tilde{t}_1 squarks are correlated with larger values of m_h (due to top-Yukawa radiative corrections to m_h)
 - in many models with $m_A \gg M_Z$, then $h \simeq H_{\text{SM}}$: the LEP2 lower bound of 114.1 GeV applicable
- $m_{\tilde{W}_1}$ vs. $m_{\tilde{\tau}_1}$
 - dashed lines: reach of ILC500 ($\sqrt{s} = 500$ GeV)
 - dotted lines: reach of ILC1000 ($\sqrt{s} = 1000$ GeV)

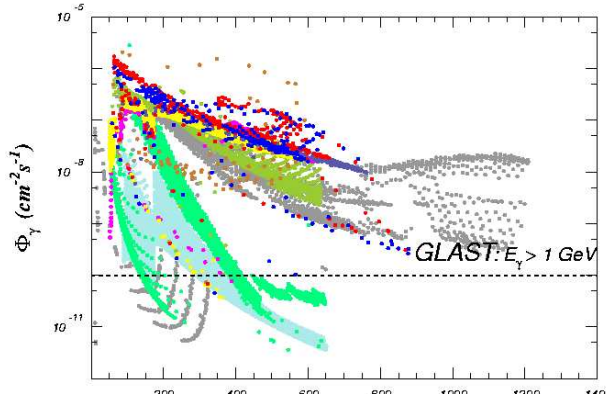
Implications for direct/indirect(neutrino) DM detection



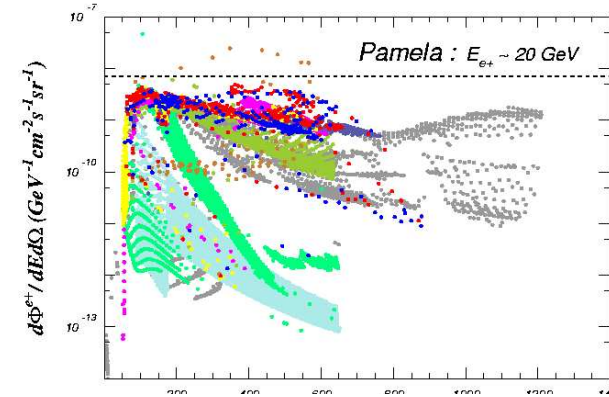
- models with **WTN** within reach of next generation of detectors
- models adjusted masses to get WMAP value below sensitivities of detectors
- muon fluxes from neutralino annihilation in the solar core to ν_{μ} states
- main contribution comes from Z -exchange ← enhanced if neutralino has high higgsino content

Implications for indirect(γ -ray, antiparticle) DM detection

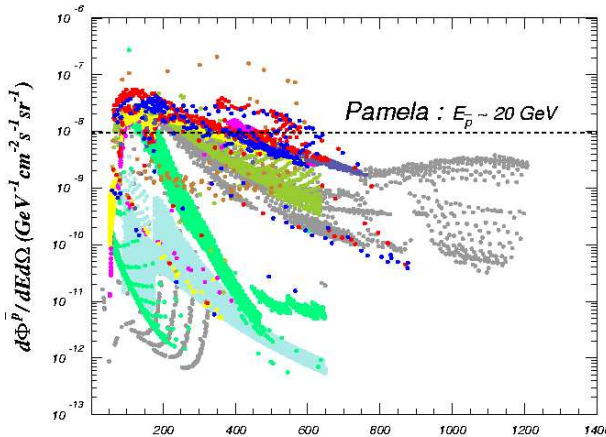
Gamma-ray Detection : Ad. Contr. N03 HM



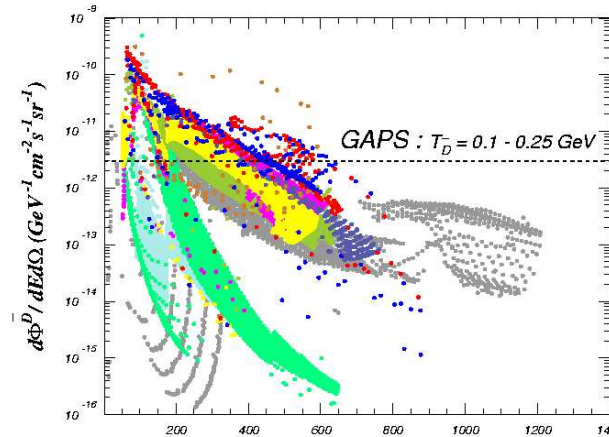
Positron Detection : Ad. Contr. N03 HM



Anti-proton Detection : Ad. Contr. N03 HM



Anti-deuteron Detection : Ad. Contr. N03 HM



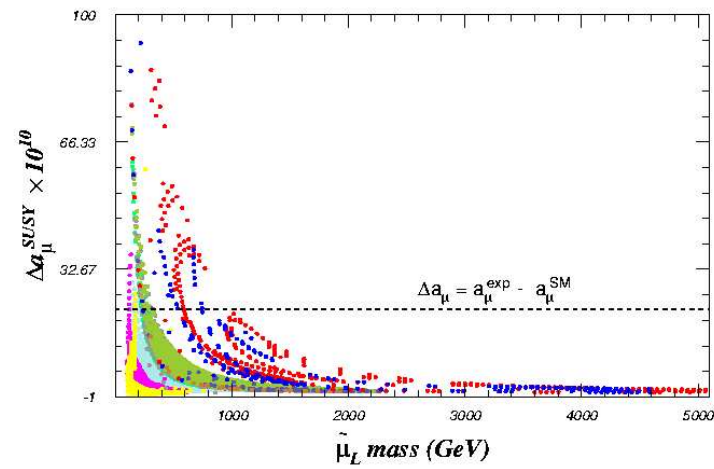
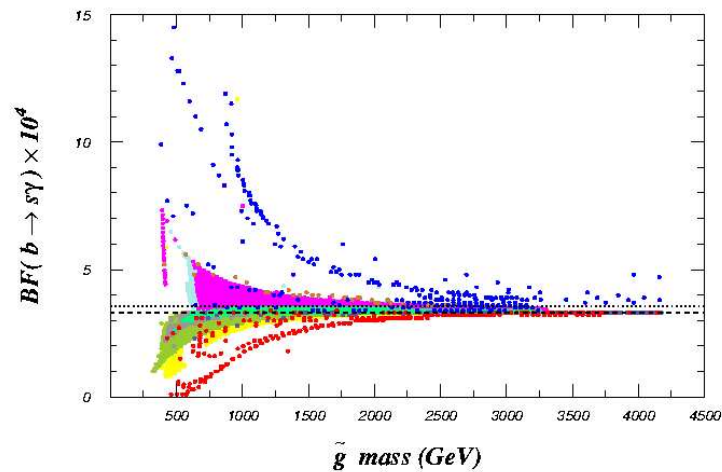
- mSUGRA : $\mu > 0$
- mSUGRA : $\mu < 0$
- NUHM1 _{μ}
- NUHM1_A
- MWDM1
- MWDM2
- BWCA2
- LM3DM
- HM2DM : $M_2 > 0$
- HM2DM : $M_2 < 0$

- mSUGRA : $\mu > 0$
- mSUGRA : $\mu < 0$
- NUHM1 _{μ}
- NUHM1_A
- MWDM1
- MWDM2
- BWCA2
- LM3DM
- HM2DM : $M_2 > 0$
- HM2DM : $M_2 < 0$

Conclusions

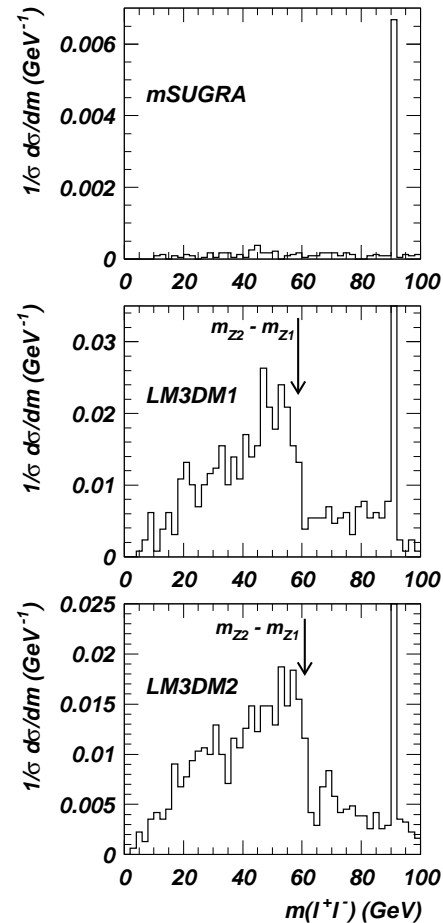
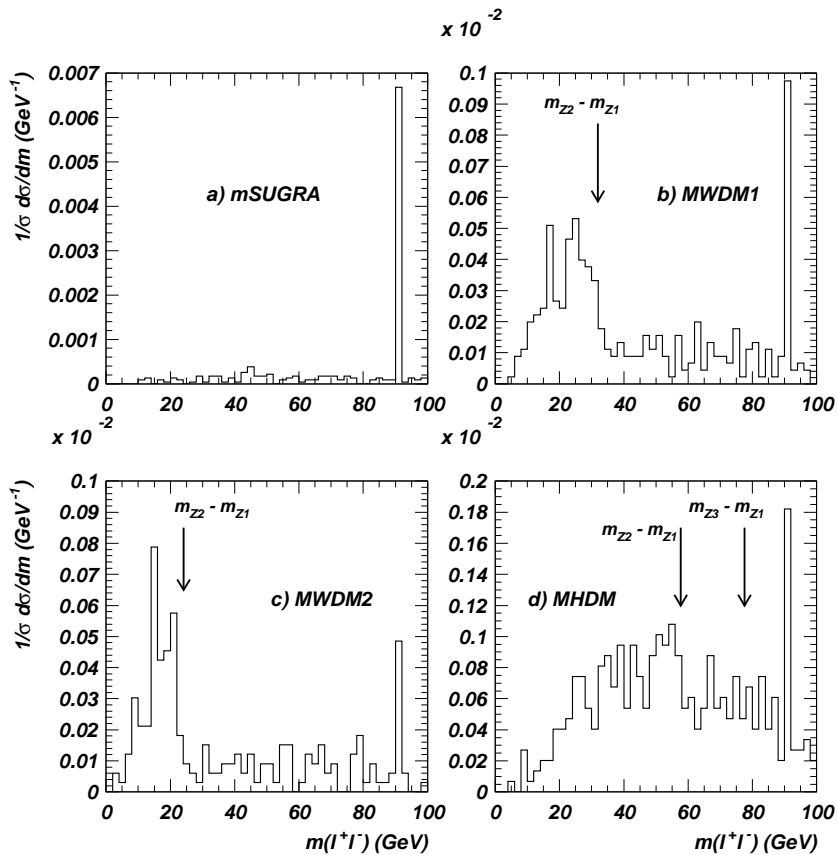
1. ★ WTN occurs *only* in FP/HB region in mSUGRA (MHDM: $m_{\tilde{q}} \gg m_{\tilde{Z}_1, \tilde{W}_1, \tilde{g}}$).
 But, in relic-density-consistent models, easily get WTN with $m_{\tilde{q}} \sim m_{\tilde{g}}$
 ★ Higgs funnel enhancement is *only* for very large $\tan\beta$ values in mSUGRA.
 But, in non-universal Higgs mass models, we have Higgs funnel for any $\tan\beta$ value
2. In many relic-density-consistent models, $\tilde{Z}_2 - \tilde{Z}_1$ mass gap $< M_Z$
 → 2-body decay modes kinematically closed
 → 3-body decay modes open \Rightarrow at least one dilepton mass edge detectable at LHC
 → location of dilepton mass edge is clean signature of SUSY models
3. ★ $m_{\tilde{q}} = m_{\tilde{g}}, m_{\tilde{q}, \tilde{g}} < 3100$ GeV for most relic-density-consistent models
 → implies SUSY signals at LHC
 ★ $m_{\tilde{\tau}} < 500$ GeV for LM3DM
 → accessible at ILC with $\sqrt{s}=1$ TeV
4. In WTN models,
 ★ enhanced annihilation rates enhance direct DM detection rates
 ★ in many cases, muon neutrino signals accessible at IceCube
 ★ indirect DM searches in galactic halo into gamma rays and anti-matter elevated; large uncertainties associated with unknown galactic DM density profile

Implications for $BF(b \rightarrow s\gamma)$ and $(g - 2)_\mu$



- mSUGRA : $\mu > 0$
 - mSUGRA : $\mu < 0$
 - NUHM1 $_\mu$
 - NUHM1 $_A$
- MWDM1
 - MWDM2
 - BWCA2
 - LM3DM
- HM2DM : $M_2 > 0$
 - HM2DM : $M_2 < 0$

Dilepton Distribution at LHC



- **mSUGRA** :
sharp peak at $m(l^+l^-) \sim M_Z$ from $\tilde{Z}_2 \rightarrow \tilde{Z}_1 Z^0$ decays
- **NUGM** :
 Z^0 peak from $\tilde{Z}_3, \tilde{Z}_4, \tilde{W}_2$ decays
+ continuum distribution
 $m(l^+l^-) < m_{\tilde{Z}_2} - m_{\tilde{Z}_1}$

MSSM RGEs

$$\frac{dm_{H_u}^2}{dt} = \frac{2}{16\pi^2} \left(-\frac{3}{5}g_1^2 M_1^2 - 3g_2^2 M_2^2 + \frac{3}{10}g_1^2 S + 3f_t^2 X_t \right)$$

$$\frac{dm_{H_d}^2}{dt} = \frac{2}{16\pi^2} \left(-\frac{3}{5}g_1^2 M_1^2 - 3g_2^2 M_2^2 - \frac{3}{10}g_1^2 S + 3f_b^2 X_b + f_\tau^2 X_\tau \right)$$

$$\frac{dm_{Q_3}^2}{dt} = \frac{2}{16\pi^2} \left(-\frac{1}{15}g_1^2 M_1^2 - 3g_2^2 M_2^2 - \frac{16}{3}g_3^2 M_3^2 + \frac{1}{10}g_1^2 S + f_t^2 X_t + f_b^2 X_b \right)$$

$$\frac{dm_{\tilde{t}_R}^2}{dt} = \frac{2}{16\pi^2} \left(-\frac{16}{15}g_1^2 M_1^2 - \frac{16}{3}g_3^2 M_3^2 - \frac{2}{5}g_1^2 S + 2f_t^2 X_t \right)$$

$$\frac{dm_{\tilde{b}_R}^2}{dt} = \frac{2}{16\pi^2} \left(-\frac{4}{15}g_1^2 M_1^2 - \frac{16}{3}g_3^2 M_3^2 + \frac{1}{5}g_1^2 S + 2f_b^2 X_b \right)$$

$$\frac{dm_{L_3}^2}{dt} = \frac{2}{16\pi^2} \left(-\frac{3}{5}g_1^2 M_1^2 - 3g_2^2 M_2^2 - \frac{3}{10}g_1^2 S + f_\tau^2 X_\tau \right)$$

$$\frac{dm_{\tilde{\tau}_R}^2}{dt} = \frac{2}{16\pi^2} \left(-\frac{12}{5}g_1^2 M_1^2 + \frac{3}{5}g_1^2 S + 2f_\tau^2 X_\tau \right)$$

$$S = m_{H_u}^2 - m_{H_d}^2 + Tr \left[\mathbf{m}_Q^2 - \mathbf{m}_L^2 - 2\mathbf{m}_U^2 + \mathbf{m}_D^2 + \mathbf{m}_E^2 \right]$$

where $t = \log(Q)$, $f_{t,b,\tau}$ are the t , b and τ Yukawa couplings, and

$$\begin{aligned} X_t &= m_{Q_3}^2 + m_{\tilde{t}_R}^2 + m_{H_u}^2 + A_t^2 \\ X_b &= m_{Q_3}^2 + m_{\tilde{b}_R}^2 + m_{H_d}^2 + A_b^2 \\ X_\tau &= m_{L_3}^2 + m_{\tilde{\tau}_R}^2 + m_{H_d}^2 + A_\tau^2 \end{aligned}$$

Some Benchmark Cases: non-universal scalar mass models

parameter	mSUGRA	NMH	NUHM1 $_{\mu}$	NUHM1 $_A$	HS
special	—	$m_0(1, 2)$	m_{ϕ}	m_{ϕ}	δ_H
value	—	54	549	-728	-1.36
μ	385.1	386.5	105.8	748.5	269.3
$m_{\tilde{g}}$	729.7	722.1	731.4	733.4	728.9
$m_{\tilde{u}_L}$	720.8	658.4	724.3	720.5	720.1
$m_{\tilde{t}_1}$	523.4	526.5	484.1	624.5	505.8
$m_{\tilde{b}_1}$	656.8	659.8	642.2	689.5	645.4
$m_{\tilde{e}_L}$	364.5	216.2	364.8	365.8	373.4
$m_{\tilde{e}_R}$	322.3	128.9	322.5	321.9	301.8
$m_{\tilde{\tau}_1}$	317.1	317.6	317.8	316.4	299.3
$m_{\tilde{W}_2}$	411.7	412.7	264.7	754.8	321.1
$m_{\tilde{W}_1}$	220.7	219.5	91.1	234.9	196.6
$m_{\tilde{Z}_2}$	220.6	219.4	117.4	234.5	198.1
$m_{\tilde{Z}_1}$	119.2	118.4	69.0	121.5	115.4
m_A	520.3	521.9	584.5	268.5	279.0
m_{H^+}	529.8	531.4	593.8	281.6	292.0
m_h	110.1	110.1	109.8	110.5	109.8
$\Omega_{\tilde{Z}_1} h^2$	1.1	0.10	0.11	0.11	0.10
$\sigma_{SI}(\tilde{Z}_1 p)$	2.1×10^{-9} pb	2.1×10^{-9} pb	7.8×10^{-8} pb	1.2×10^{-9} pb	2.7×10^{-8} pb
$R_{\tilde{H}}$	0.15	0.14	0.84	0.06	0.26

Some Benchmark Cases: non-universal gaugino mass models

parameter	mSUGRA	MWDM	BWCA	LM3DM	HM2DM
special	—	$M_1(M_{GUT})$	$M_1(M_{GUT})$	$M_3(M_{GUT})$	$M_2(M_{GUT})$
value	—	490	-480	160	900
μ	385.1	385.9	376.6	185.3	134.8
$m_{\tilde{g}}$	729.7	729.9	731.7	420.2	736.4
$m_{\tilde{u}_L}$	720.8	721.2	722.0	496.9	901.8
$m_{\tilde{u}_R}$	702.7	708.9	709.9	467.0	696.3
$m_{\tilde{t}_1}$	523.4	526.5	536.3	312.2	394.3
$m_{\tilde{b}_1}$	656.8	656.0	658.9	443.2	686.4
$m_{\tilde{e}_L}$	364.5	371.5	371.4	366.1	669.3
$m_{\tilde{e}_R}$	322.3	353.3	352.2	322.6	321.3
$m_{\tilde{W}_2}$	411.7	412.4	404.5	282.9	719.7
$m_{\tilde{W}_1}$	220.7	220.8	220.0	152.5	136.5
$m_{\tilde{Z}_2}$	220.6	223.2	219.2	163.6	142.3
$m_{\tilde{Z}_1}$	119.2	194.6	201.7	105.5	94.8
m_A	520.3	525.9	518.6	398.3	670.7
m_{H^+}	529.8	535.3	528.1	408.7	679.8
m_h	110.1	110.2	109.8	106.0	111.9
$\Omega_{\tilde{Z}_1} h^2$	1.1	0.10	0.10	0.10	0.10
$\sigma_{SI}(\tilde{Z}_1 p)$	2.1×10^{-9} pb	1.5×10^{-8} pb	3.1×10^{-11} pb	7.2×10^{-8} pb	3.4×10^{-8} pb
$R_{\tilde{H}}$	0.15	0.25	0.16	0.50	0.67